# Small Quantum Cohomology as Quantum Integrable System

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#### Outline

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- Reminder: (small) quantum cohomology
- Quantum Kostka numbers and toric Schur functions
- Vicious and osculating walkers on the cylinder
- Free fermion description

Related literature

### Related literature (http://www.maths.gla.ac.uk/~ck/publications.html)

#### quantum cohomology/WZNW fusion rings as integrable models:

- C. K. and C. Stroppel, Adv Math 225 (2010) 200-68; arXiv:0909.2347 (type A: g = û(n)<sub>k</sub> and sû(n)<sub>k</sub>)
- C. K. A combinatorial derivation of the Racah-Speiser algorithm for Gromov-Witten invariants; arXiv:0910.3395
- C. K. J Phys A 43 (2010) 434021; arxiv:1006.4710
- C. K. QC via vicious and osculating walkers; arxiv:1204.4109

#### Bethe vectors as idempotents: discrete quantum NLS model

- C. K. RIMS Kokyuroku Bessatsu B28 (2011) pp. 121-53; arxiv:1106.5342 (Section 5, Cor 5.3 and Section 7)
- C. K. Cylindric Macdonald functions and a deformation of the Verlinde algebra. preprint; arxiv:1110.6356 (Section 7, Prop 7.15)

L The small quantum cohomology ring

 $\operatorname{Gr}_{n,n+k}$  Grassmannian of *n*-planes in  $\mathbb{C}^{n+k}$ 

• small quantum cohomology ring [Siebert-Tian 1997]

 $qH^*(\mathrm{Gr}_{n,n+k}) \cong \mathbb{Z}[q][e_1,\ldots,e_n]/\langle h_{k+1},\ldots,h_{n+k-1},h_{n+k}+q(-1)^n\rangle$ 

where  $h_r = \det(e_{1-i+j})_{1 \le i,j \le r}$  and a vector space basis is given by  $\{s_{\lambda} := \det(e_{\lambda'_i - i+j})_{1 \le i,j \le k}\}$  with

 $\lambda \in \{ \text{partitions with Young diagram in } n \times k \text{ box } \}$ 

 Fusion ring of û(n)<sub>k</sub> Wess-Zumino-Novikov-Witten model [Gepner, Intriligator, Vafa, Witten] and [Agnihotri]:

$$\mathcal{F}_{n,k}^{\mathbb{Z}} \cong qH^*(\mathrm{Gr}_{n,n+k})/\langle q-1 \rangle$$

 $\mathcal{F}_{n,k} := \mathcal{F}_{n,k}^{\mathbb{Z}} \otimes_{\mathbb{Z}} \mathbb{C}$  is called *Verlinde algebra*.

Reminder

### Schubert varieties and rational maps

Given a flag  $F_1 \subset F_2 \subset \cdots \subset F_{n+k} = \mathbb{C}^{n+k}$  the Schubert variety  $\Omega_{\lambda}(F)$  is defined as

$$\Omega_{\lambda}(F) = \{ V \in \operatorname{Gr}_{n,n+k} \mid \dim(V \cap F_{k+i-\lambda_i}) \geq i, i = 1, \dots n \}.$$

#### Definition of 3-point Gromov-Witten invariants

 $C_{\lambda,\mu}^{\nu,d} = \# \text{ of rational } f : \mathbb{P}^1 \to \operatorname{Gr}_{n,N} \text{ of degree } d \text{ which meet}$  $\Omega_{\lambda}(F), \ \Omega_{\mu}(F'), \ \Omega_{\nu^{\vee}}(F'') \text{ for general flags } F, F', F'' \text{ modulo}$ automorphisms in  $\mathbb{P}^1$ . If there is an  $\infty$  number of such maps, set  $C_{\lambda,\mu}^{\nu,d} = 0.$ 

Poincaré duality:  $\nu^{\vee} = (k - \nu_n, \dots, k - \nu_1)$ Schubert class:  $[\Omega_{\lambda}] \mapsto s_{\lambda}$ 

Quantum Kostka numbers and Gromov-Witten invariants

# Quantum Kostka numbers and Gromov-Witten invariants

[Bertram, Ciocan-Fontanine, Fulton]:

$$s_{\mu} \star s_{\lambda_{1}} \star \cdots \star s_{\lambda_{r}} = \sum_{d \ge 0, \nu \in (n,k)} q^{d} s_{\nu} K_{\nu/d/\mu,\lambda}$$
$$s_{\mu} \star s_{(1^{\lambda_{1}})} \star \cdots \star s_{(1^{\lambda_{r}})} = \sum_{d \ge 0, \nu \in (n,k)} q^{d} s_{\nu} K_{\nu'/d/\mu',\lambda}$$

Quantum Giambelli formula [Bertram]:  $s_{\lambda} = \det(s_{\lambda_i - i + j})$ 

$$s_{\mu} \star s_{\lambda} = \sum_{d \ge 0, \nu \in (n,k)} q^d C_{\lambda\mu}^{\nu,d} s_{\nu}, \qquad d = \frac{|\lambda| + |\mu| - |\nu|}{n+k}$$

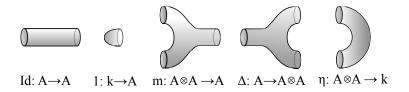
3-point, genus 0 Gromov-Witten invariants

Quantum Kostka numbers and Gromov-Witten invariants

#### Proposition (intersection pairing)

 $\mathcal{F}_{n,k}$  is a commutative Frobenius algebra with  $\eta(s_{\lambda}, s_{\mu}) = \delta_{\lambda^{\vee}\mu}$ .

A Frobenius algebra A is a finite-dimn'l, unital, assoc algebra with non-degenerate bilinear form  $\eta(a \star b, c) = \eta(a, b \star c), a, b, c \in A$ .



#### Topological quantum field theories [Witten, Segal, Atiyah]

Commutative Frobenius algebras are categorically equivalent to 2D topological quantum field theories. [Dijkgraaf]

Introduction

Frobenius structures and toric Schur functions

### Computation of the coproduct

Frobenius isomorphism  $\Phi : s_{\lambda} \mapsto \eta(s_{\lambda}, \circ)$ 

$$\begin{array}{cccc} \mathcal{F}_{n,k} & \stackrel{\Delta}{\longrightarrow} & \mathcal{F}_{n,k} \otimes \mathcal{F}_{n,k} \\ & & \downarrow \Phi & & \downarrow \Phi \otimes \Phi \\ \mathcal{F}_{n,k}^* & \stackrel{\mathfrak{m}^*}{\longrightarrow} & \mathcal{F}_{n,k}^* \otimes \mathcal{F}_{n,k}^* \end{array}$$

Proposition (generalised skew Schur function)

$$\Delta s_
u = \sum_{d,\mu} s_{
u/d/\mu} \otimes s_\mu, \quad s_{
u/d/\mu} := \sum_\lambda C^{
u,d}_{\lambda\mu} s_\lambda \;.$$

Introduction

Frobenius structures and toric Schur functions

Toric Schur function [Postnikov]:

$$s_{\nu/d/\mu}(x_1,\ldots,x_n) = \sum_{\lambda} C_{\lambda\mu}^{\nu,d} s_{\lambda}(x_1,\ldots,x_n)$$
  
= 
$$\sum_{\lambda} K_{\nu/d/\mu,\lambda} m_{\lambda}(x_1,\ldots,x_n),$$

∞-many variables: cylindric Schur functions [Gessel-Krattenthaler] [McNamara] [Lapointe-Morse] [Lam]

Fusion ring as quantum integrable model (Korff-Stroppel 2010)

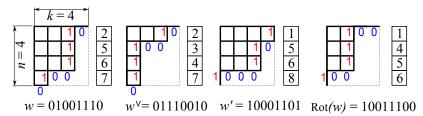
Identify the toric Schur functions as partition functions and the fusion ring as the quantum integrals of motion.

Other example: XXX Bethe algebra [Varchenko et al]

└─Vicious and osculating walkers

01-words and partitions

Reminder: correspondence between 01-words and partitions



The following bijections induce symmetries of GW invariants:

$$w \mapsto w^{\vee} = w_N \dots w_2 w_1$$
  

$$w \mapsto w' = (1 - w_N) \cdots (1 - w_2)(1 - w_1)$$
  

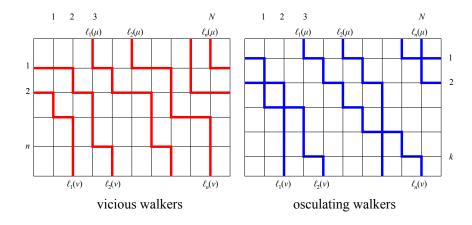
$$w \mapsto \operatorname{Rot}(w) := w_2 w_3 \dots w_N w_1$$

Position of 1-letters:  $\ell_i(\lambda) = \lambda_{n+1-i} + i$ 

└─Non-intersecting lattice paths

### Vicious and osculating walkers on the cylinder

Statistical models on  $n \times N$  and  $k \times N$  square lattice with periodic boundary conditions in the horizontal direction (N = n + k).

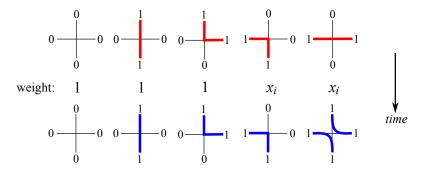


Vicious and osculating walkers

└─Non-intersecting lattice paths

#### Allowed vertex configurations and their weights

 $x_i$  indeterminate assigned to the  $i^{th}$  lattice row.



(percolation c.f. [Brak][Fisher][Forrester][Guttmann et al][Wu])

transfer matrices

### Row partition functions

Fix start/end positions via 01-words  $w(\mu)$ ,  $w(\nu)$  of length N.

#### Definition (transfer matrices)

Weighted sums over row configurations:

$$\begin{split} E(x_i)_{\nu,\mu} &:= \sum_{osc \ row \ config} q^{\frac{\# \ of \ outer \ edges}{2}} x_i^{\# \ of \ horizontal \ edges} \\ H(x_i)_{\nu,\mu} &:= \sum_{vicious \ row \ config} q^{\frac{\# \ of \ outer \ edges}{2}} x_i^{\# \ of \ horizontal \ edges} \end{split}$$

Proposition (integrability  $\equiv$  commuting transfer matrices)  $E(x)E(y) = E(y)E(x), \ H(x)H(y) = H(y)H(x), \ E(x)H(y) = H(y)E(x)$ 

Weighted path counting

Theorem (Generating function for Gromov-Witten invariants)

The partition functions have the following expansions,

$$(H(x_n)\cdots H(x_2)\cdot H(x_1))_{\nu,\mu} = \sum_{d\geq 0} q^d s_{\nu/d/\mu}(x_1,\ldots,x_n)$$
  
$$(E(x_k)\cdots E(x_2)\cdot E(x_1))_{\nu,\mu} = \sum_{d\geq 0} q^d s_{\nu'/d/\mu'}(x_1,\ldots,x_k)$$

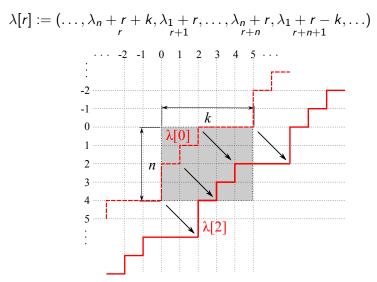
Let h(s), c(s) be hook length and content of  $s \in \lambda$ . Corollary (Sum rule for Gromov-Witten invariants)

Set  $x_i = q = 1$  for all  $1 \le i \le n$ . Then

$$H_{\nu,\mu}^{n} = \sum_{d,\lambda} C_{\lambda\mu}^{\nu,d} \prod_{s \in \lambda} \frac{n + c(s)}{h(s)} = E_{\nu',\mu'}^{n}$$

Proof

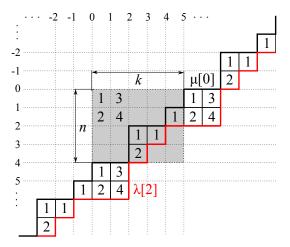
### Cylindric loops



Vicious and osculating walkers

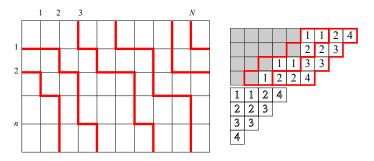
### Cylindric skew tableaux

$$\lambda/d/\mu := \{ \langle i,j \rangle \in \mathbb{Z} \times \mathbb{Z}/(n,-k)\mathbb{Z} \mid \lambda[d]_i \ge j > \mu[0]_i \}$$



#### Proposition

Vicious/osculating paths are in bijection with cylindric tableaux.



Level-rank duality:  $\tau \circ H = E \circ \tau$  with  $\tau : \lambda \mapsto \lambda'$ 

 $\mathcal{K}_{
u/d/\mu,\lambda}=\#$  of cylindric tableaux of weight  $\lambda$  [BCF][Postnikov]

Vicious and osculating walkers

Quantum integrals of motion: XX spin-chain

### Quantum integrals of motion

Define matrices  $S_{(a|b)}$  via the expansion

$$H(x)E(y) = 1 + (x + y) \sum_{a,b \ge 0} x^{a}y^{b}S_{(a|b)}$$

#### Definition (Fusion matrices)

Let 
$$\lambda = (\alpha_1, \dots, \alpha_r | \beta_1, \dots, \beta_r)$$
 with  $\lambda_1, \lambda'_1 < N$   
 $S_\lambda := \det(S_{(\alpha_i | \beta_i)})_{1 \le i, j \le r}$ 

Proposition (Functional relation)  $H(x)E(-x) = 1 + (-1)^n q x^{n+k}$ 

└─ The algebraic Bethe ansatz

#### Theorem (Korff-Stroppel 2010)

There exists an orthogonal basis  $\{\mathfrak{e}_{\lambda}\}_{\lambda \in (n,k)}$  such that

- **1** the matrices H, E and the  $S_{\lambda}$ 's are diagonal.
- **2** mapping the  $\mathfrak{e}_{\lambda}$ 's onto the idempotents of  $\mathcal{F}_{n,k}$  yields an algebra isomorphism, in particular

$$\mathcal{S}_\lambda \mathcal{S}_\mu = \sum_{d \geq 0, 
u \in (n,d)} q^d \, \mathcal{C}_{\lambda\mu}^{
u,k} \mathcal{S}_
u \; .$$

#### XX-Heisenberg spin chain

The transfer matrices H, E commute with the Hamiltonian of the so-called quantum XX-Heisenberg spin chain.

Free fermion description

Clifford algebra

### Fermion creation and annihilation

Fix N = n + k and consider the vector space (Fock space)

$$\mathcal{F} = \bigoplus_{n=0}^{N} \mathcal{F}_{n,k}, \qquad \mathcal{F}_{n,k} = \mathbb{C} W_{n,k},$$

where  $\mathcal{F}_{0,N} = \mathbb{C}\{0\cdots 0\} = \mathbb{C}$  and  $w = 0\cdots 0$  is the vacuum  $\varnothing$ . Let  $n_i(w) = w_1 + \cdots + w_i$  be the number of 1-letters in [1, i]. For  $1 \le i \le N$  define the (linear) maps  $\psi_i^*, \psi_i : \mathcal{F}_{n,k} \to \mathcal{F}_{n\pm 1,k\mp 1}$ ,

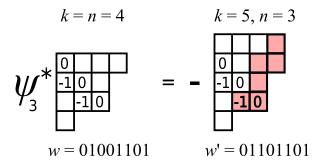
$$\begin{split} \psi_i^*(w) &:= \begin{cases} (-1)^{n_{i-1}(w)} w', & w_i = 0 \text{ and } w'_j = w_j + \delta_{i,j} \\ 0, & w_i = 1 \end{cases} \\ \psi_i(w) &:= \begin{cases} (-1)^{n_{i-1}(w)} w', & w_i = 1 \text{ and } w'_j = w_j - \delta_{i,j} \\ 0, & w_i = 0. \end{cases} \end{split}$$

-Free fermion description

-Clifford algebra

#### Example

Take 
$$n = k = 4$$
 and  $\mu = (4, 3, 3, 1)$ .



The boundary ribbon (shaded boxes) starts in the (3 - n) = -1 diagonal. Below the diagram the respective 01-words  $w(\mu)$  and  $w(\psi_3^*\mu)$  are displayed.

-Free fermion description

Clifford algebra

#### Proposition (Clifford algebra)

The maps  $\psi_i, \psi_i^* : \mathcal{F}_n \to \mathcal{F}_{n \mp 1, k \pm 1}$  yield an irred rep of the Clifford algebra, i.e. one has the relations (i, j = 1, ..., N)

$$\psi_i \psi_j + \psi_j \psi_i = \psi_i^* \psi_j^* + \psi_j^* \psi_i^* = 0, \qquad \psi_i \psi_j^* + \psi_j^* \psi_i = \delta_{ij}.$$

Introducing  $\langle w, w' \rangle = \prod_i \delta_{w_i, w'_i}$  one has  $\langle \psi_i^* w, w' \rangle = \langle w, \psi_i w' \rangle$  for any  $w, w' \in \mathcal{F}$ .

Bijections: partitions - 01-words - words in the Clifford algebra

$$\lambda \longleftrightarrow w(\lambda) \longleftrightarrow \psi^*_{\ell_1(\lambda)} \cdots \psi^*_{\ell_n(\lambda)} \varnothing$$

The Clifford algebra is the fundamental object in the description of quantum cohomology.

Free fermion description

Return of the vicious walkers

### Connection with vicious walkers

Expansion of the transfer matrix  $H(x) : \mathcal{F}_{n,k} \to \mathcal{F}_{n,k}$ ,

$$H(x) = \sum_{r=0}^{N} x^{r} H_{r}, \quad H_{r} = \sum_{|\alpha|=r} \psi_{N}^{\alpha_{N}} u_{N-1}^{\alpha_{N}-1} \cdots u_{1}^{\alpha_{1}} ((-1)^{r-1} q \psi_{1}^{*})^{\alpha_{N}}$$

where  $u_i = \psi_{i+1}^* \psi_i$  shifts one particle from site *i* to site i + 1.

Proposition (commutation relation with fusion matrices)

$$S_{\lambda}(q)\psi_i^* = \psi_i^*S_{\lambda}(-q) + \sum_{r=1}^{\ell(\lambda)}\psi_{i+r}^*\sum_{\lambda/\mu=(r)}S_{\mu}(-q)$$

where  $\psi_{j+N}^* = (-1)^{n+1} q \psi_j^*$  and n = particle number operator.

-Free fermion description

Fermion creation of quantum cohomology rings

# Fermion creation of quantum cohomology rings Corollary (Korff-Stroppel 2010)

The last commutation relation implies the product formula

$$\lambda \star \psi_i^*(\mu) = S_{\lambda}(q)\psi_i^*(\mu) = \sum_{r=0}^{\lambda_1} \sum_{\lambda/\nu=(r)} \psi_{i+r}^*(\nu \neq \mu)$$

where  $\bar{\star}$  denotes the product with q replaced by -q.

#### Inductive algorithm

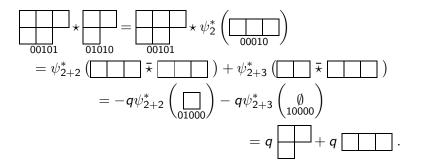
One can successively generate the entire ring hierarchy  $\{qH^*(Gr_{n,n+k})\}_{n=0}^{n+k}$  starting from n = 0.

-Free fermion description

Fermion creation of quantum cohomology rings

#### Example

Consider the ring  $qH^*(\operatorname{Gr}_{2,5})$ . Via  $\psi_i^*: qH^*(\operatorname{Gr}_{1,5}) \to qH^*(\operatorname{Gr}_{2,5})$ one can compute the product in  $qH^*(\operatorname{Gr}_{2,5})$  through the product in  $qH^*(\operatorname{Gr}_{1,5})$ :



Free fermion description

Fermion creation of quantum cohomology rings

### Fermionic product formula

Let  $\lambda, \mu \in (n, k)$ . Then

$$w(\lambda)\star w(\mu) := \sum_{T} \psi^*_{\ell_1(\mu)+t_n} \overline{\psi}^*_{\ell_2(\mu)+t_{n-1}} \psi^*_{\ell_3(\mu)+t_{n-2}} \overline{\psi}^*_{\ell_4(\mu)+t_{n-3}} \cdots \varnothing,$$

where

- $\ell_i(\mu)$  positions of 1-letters in  $w(\mu)$
- $T = (\text{semistandard}) \text{ tableau of shape } \lambda$
- $t_i$  = number of entries  $1 \le i \le n$  in T

• 
$$ar{\psi}^*_i = \psi^*_i$$
 for  $i=1,\ldots,N$  and

$$\psi_{i+N}^* := (-1)^{n+1} q \psi_i^*, \quad \bar{\psi}_{i+N}^* := (-1)^n q \bar{\psi}_i^*, \quad n := \sum_{i=1}^N \psi_i^* \psi_i.$$

Free fermion description

Fermion creation of quantum cohomology rings

### Example

Set 
$$N = 7$$
,  $n = N - k = 4$  and  $\lambda = (2, 2, 1, 0)$ ,  $\mu = (3, 3, 2, 1)$ .

Step 1. Positions of 1-letters:  $\ell(\mu) = (\ell_1, \dots, \ell_4) = (2, 4, 6, 7)$ . Step 2. Write down all tableaux of shape  $\lambda$  such that  $\ell' = (\ell_1 + t_n, \dots, \ell_n + t_1) \mod N$  with  $\ell'_i \neq \ell'_j$  for  $i \neq j$ .

$$\begin{array}{c} 11\\ 24\\ 3\\ 3\\ (3,5,7,9) \end{array}, \begin{array}{c} 11\\ 223\\ 4\\ 4\\ \end{array}, \begin{array}{c} 11\\ 222\\ 4\\ 4\\ \end{array}, \begin{array}{c} 11\\ 34\\ 4\\ 4\\ \end{array}, \begin{array}{c} 13\\ 24\\ 3\\ 3\\ \end{array}, \begin{array}{c} 12\\ 33\\ 4\\ 4\\ \end{array}, \begin{array}{c} 22\\ 33\\ 4\\ 4\\ \end{array}, \begin{array}{c} 12\\ 34\\ 4\\ 4\\ \end{array}, \begin{array}{c} 22\\ 34\\ 4\\ 4\\ \end{array}, \begin{array}{c} 34\\ 4\\ 4\\ \end{array}, \begin{array}{c} 22\\ 34\\ 4\\ 4\\ \end{array}, \begin{array}{c} 34\\ 4\\ 4\\ \end{array}, \begin{array}{c} 34\\ 4\\ 4\\ \end{array}, \begin{array}{c} 36\\ 36\\ 36\\ 78\\ \end{array}, \begin{array}{c} 36, 78\\ 36, 78\\ (3,6,8,7) \end{array}, \begin{array}{c} 36, 8,7\\ (4,5,7,8) \end{array}, \begin{array}{c} 45, 7,8\\ (4,5,7,8) \end{array}, \begin{array}{c} 45, 8,7\\ (4,5,7,8) \end{array}, \begin{array}{c} 45, 8,7\\ (4,5,8,7) \end{array}$$

Step 3. For each  $\ell'_i > N$  make the replacement

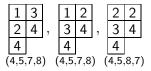
$$\psi_{\ell'_1}^*\cdots\psi_{\ell'_n}^*\varnothing\to (-1)^{n+1}q\psi_{\ell'_1}^*\cdots\psi_{\ell'_n-N}^*\cdots\psi_{\ell'_n}^*\varnothing$$

Step 4. Let  $\ell''$  be the reduced positions in [1, N]. Choose permutation  $\pi \in S_n$  s.t.  $\ell''_1 < \cdots < \ell''_n$  and multiply with  $(-1)^{\ell(\pi)}$ .

Free fermion description

Fermion creation of quantum cohomology rings

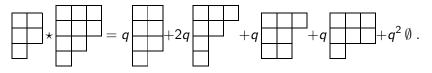
The three tableaux



yield the same 01-word w = 1001101,  $\lambda(w) = (3, 2, 2, 0)$  but with changing sign,

$$\begin{split} \psi_{\ell_1+2}^*\bar{\psi}_{\ell_2+1}^*\psi_{\ell_3+1}^*\bar{\psi}_{\ell_4+1}^*\varnothing &= \psi_{\ell_1+2}^*\bar{\psi}_{\ell_2+1}^*\psi_{\ell_3+1}^*\bar{\psi}_{\ell_4+1}^*\varnothing &= \\ &-\psi_{\ell_1+2}^*\bar{\psi}_{\ell_2+1}^*\psi_{\ell_3+2}^*\bar{\psi}_{\ell_4}^*\varnothing &= q\;\psi_1^*\psi_4^*\psi_5^*\psi_7^*\varnothing \;. \end{split}$$

We obtain the product expansion



Free fermion description

Fermion creation of quantum cohomology rings

#### Corollary (Quantum Racah-Speiser Algorithm, C.K. 2009)

Let  $\lambda, \mu, \nu \in \mathfrak{P}_{n,k}$ . Given a permutation  $\pi \in S_n$  set

$$\begin{aligned} \alpha_i(\pi) &= (\ell_i(\nu) - \ell_{\pi(i)}(\mu)) \mod \mathsf{N} \ge \mathsf{0} \\ d(\pi) &= \#\{i \mid \ell_i(\nu) - \ell_{\pi(i)}(\mu) < \mathsf{0}\} \;. \end{aligned}$$

Then one has the following identity for Gromov-Witten invariants,

$$\mathcal{C}_{\lambda\mu}^{
u,d} = \sum_{\pi\in \mathcal{S}_n,\,d(\pi)=d} (-1)^{\ell(\pi)+(n-1)d} \,\mathcal{K}_{\lambda,lpha(\pi)} \ ,$$

where  $K_{\lambda\mu}$  are the Kostka numbers.

Setting q = 0 the formula specializes to the known Racah-Speiser algorithm for Littlewood-Richardson coefficients.

Free fermion description

L The End

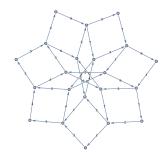
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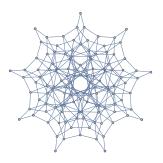
Kirillov-Reshetikhin crystals

Basic notions of crystal theory

#### Example

KR crystal:  $\mathfrak{g} = \widehat{\mathfrak{sl}}_N$ ,  $B_r = \{ \text{ 01-words with } r \text{ 1-letters} \}$ 



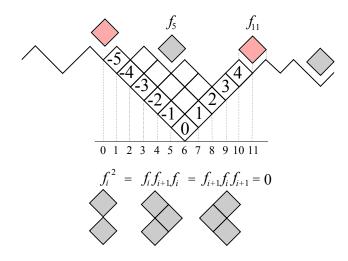


Kirillov-Reshetikhin crystals

Basic notions of crystal theory

### Example (cont'd)

#### Affine nil Temperley-Lieb algebra



Kirillov-Reshetikhin crystals

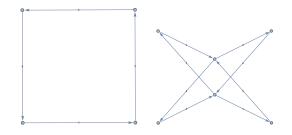
Basic notions of crystal theory

### Tensor products of crystals

 $B\otimes B'$  is the set B imes B' together with the maps,

$$egin{aligned} e_i(b\otimes b') &= egin{aligned} e_i(b)\otimes b', & arepsilon_i(b) > arphi_i(b') \ b\otimes e_i(b'), & ext{else} \end{aligned}$$
 $f_i(b\otimes b') &= egin{aligned} f_i(b)\otimes b', & arepsilon_i(b) \ge arphi_i(b') \ b\otimes f_i(b'), & ext{else} \end{aligned}$ 

where one sets  $b \otimes \emptyset = \emptyset$  and  $\emptyset \otimes b' = \emptyset$ .

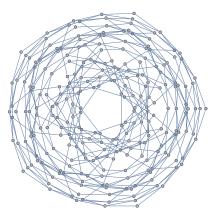


Kirillov-Reshetikhin crystals

Basic notions of crystal theory

### Things can get more complicated ....

Set N = 5 and consider the KR crystal  $B = B_2 \otimes B_1 \otimes B_1$ :



KR crystals are *perfect* crystals: they and their tensor products are always connected.

Kirillov-Reshetikhin crystals

Basic notions of crystal theory

### The combinatorial R-matrix

#### Theorem (Kashiwara et al)

There exists a unique graph isomorphism  $R_{r,s} : B_r \otimes B_s \to B_s \otimes B_r$ which preserves the crystal structure.

#### Example (c.f. Nakayashiki-Yamada)

Let N = 6 and r = 3, s = 2. Then we find

$$R_{3,2}\left(\begin{array}{c} \boxed{2}\\ 4\\ 6\end{array} \otimes \begin{array}{c} 1\\ 3\end{array}\right) = \begin{array}{c} \boxed{2}\\ 6\end{array} \otimes \begin{array}{c} 1\\ 3\\ 4\end{array}$$





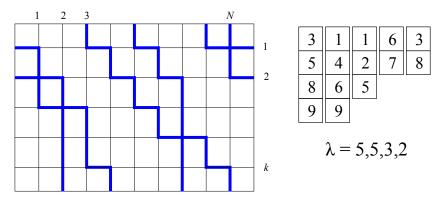
Kirillov-Reshetikhin crystals

Lattice paths as crystal vertices

### Lattice paths as crystal vertices

 $\mathcal{T}_{\nu',\mu'}(\lambda')$  cylindric tableaux of shape  $\nu'/d/\mu'$  and weight  $\lambda'$ .

Define  $\iota: \mathcal{T}_{\nu',\mu'}(\lambda') \to B_{\lambda'_1} \otimes B_{\lambda'_2} \cdots \otimes B_{\lambda'_k}$  as follows:



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-Main theorem

For simplicity assume 
$$d > d_{min} := \max_\ell \{\sum_{i=1}^\ell (w_i(
u) - w_i(\mu))\}.$$

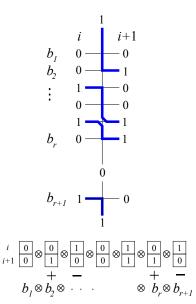
Theorem (Osculating walkers as crystal vertices)

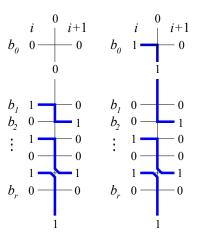
Let  $b \in B_{\lambda} := B_{\lambda'_{1}} \otimes \cdots \otimes B_{\lambda'_{k}}$ . The following are equivalent. (i)  $b \in \iota(\mathcal{T}_{\nu',\mu'}(\lambda'))$ (ii)  $\varphi(b) = \sum_{i \in I} (1 - w_{i+1}(\mu)) \omega_{i}, \quad \varepsilon(b) = \sum_{i \in I} (1 - w_{i+1}(\nu)) \omega_{i}.$ (iii)  $R_{\lambda}(b_{\mu} \otimes \operatorname{Rot}^{-1}b) = b \otimes b_{\nu}$ , where Rot is the  $\widehat{\mathfrak{sl}}_{N}$  Dynkin diagram automorphism.

Here  $R_{\lambda} := R_{n,\lambda'_r} \cdots R_{n,\lambda'_2} R_{n,\lambda'_1}$  is the unique crystal graph isomorphism  $R_{\lambda} : B_n \otimes B_{\lambda} \to B_{\lambda} \otimes B_n$ .

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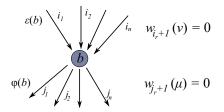
Proof: (i)  $\Rightarrow$  (ii)





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Proof: (ii)  $\Rightarrow$  (iii)



#### Claim

#### It follows from (ii) that

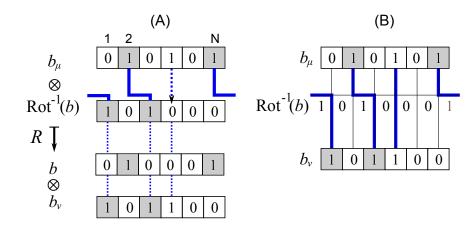
$$arepsilon(b_\mu\otimes \operatorname{Rot}^{-1}b)=arepsilon(b\otimes b_
u) ext{ and } arphi(b_\mu\otimes \operatorname{Rot}^{-1}b)=arphi(b\otimes b_
u) ext{ .}$$

Claim + uniqueness of combinatorial  $R \Rightarrow$  (iii)

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-Proof: (iii)  $\Rightarrow$  (i)

### Osculating walks via the combinatorial R-matrix



Kirillov-Reshetikhin crystals

Littlewood-Richardson coefficients

### The case of minimal degree

$$d = d_{min} := \max_{\ell} \{ \sum_{i=1}^{\ell} (w_i(\nu) - w_i(\mu)) \}.$$

Proposition (Fulton-Woodward, Yong, Postnikov)

There exists  $1 \leq a \leq N$  such that

$$K_{
u/d/\mu,\lambda} = K_{\mathrm{Rot}^{\mathfrak{s}}(
u)/\mathrm{Rot}^{\mathfrak{s}}(\mu),\lambda}$$
 and  $C_{\lambda,\mu}^{
u,d} = C_{\lambda\mathrm{Rot}^{\mathfrak{s}}(\mu)}^{\mathrm{Rot}^{\mathfrak{s}}(
u),0}$ 

Robinson-Schensted-Knuth correspondence:

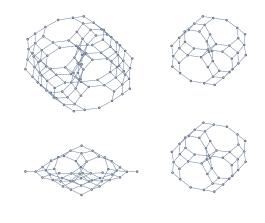
$$B_{\lambda} \cong \bigoplus_{\alpha \leq \lambda} B(\alpha) \times \mathrm{SST}(\alpha', \lambda') ,$$

where  $B(\alpha)$  is the irred  $\mathfrak{sl}_N$ -crystal of highest weight  $\alpha$ .

Kirillov-Reshetikhin crystals

Littlewood-Richardson coefficients

Previous example with N = 5 and  $B_{2,1,1} = B_2 \otimes B_1 \otimes B_1$ :



http://demonstrations.wolfram.com/KirillovReshetikhinCrystals/

Kirillov-Reshetikhin crystals

-Vicious walks and the combinatorial R-matrix

### Vicious walks

 $\tau: B_n \to B_{N-n}, \ b_{\lambda} \mapsto b_{\lambda'}$  swaps zero and one-letters in 01-word b and then reverses its order.

Define 
$$R'_{r,s} := (1 \otimes \tau) R_{N-r,s}(\tau \otimes 1)$$
 and  $R'_{\lambda} := R'_{n,\lambda_r} \cdots R'_{n,\lambda_1}$ .

#### Corollary

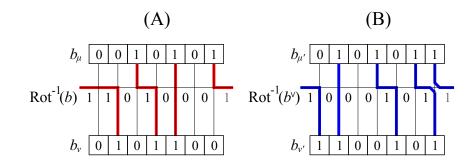
Let  $b \in B_{\lambda'}$ . The following statements are equivalent.

(i) 
$$b \in \iota(\mathcal{T}_{\nu,\mu}(\lambda))$$
  
(ii)  $\varphi(b) = \sum_{i \in I} w_{i+1}(\mu^{\vee})\omega_i, \ \varepsilon(b) = \sum_{i \in I} w_{i+1}(\nu^{\vee})\omega_i$   
(iii)  $R'_{\lambda}(b_{\mu^{\vee}} \otimes \operatorname{Rot}^{-1}b) = b \otimes b_{\nu^{\vee}}$ 

Kirillov-Reshetikhin crystals

-Vicious walks and the combinatorial R-matrix

### Constructing vicious from osculating walks



Kirillov-Reshetikhin crystals

Symmetries of quantum Kostka numbers

#### Corollary

We have the following symmetries of quantum Kostka numbers,

$$K_{\nu/d/\mu,\lambda} = K_{\nu/d/\mu,s_i\lambda} = K_{\text{Rot}(\nu)/d^R/\text{Rot}(\mu),\lambda} = K_{\mu^{\vee}/d/\nu^{\vee},\lambda},$$
  
where  $s_i\lambda = (\dots, \lambda_{i+1}, \lambda_i, \dots)$  and  $d^R = d + w_1(\mu) - w_1(\nu).$ 

Proof

- Yang-Baxter equation:  $R_{23} \circ R_{12} \circ R_{23} = R_{12} \circ R_{23} \circ R_{12}$
- Rotation (Dynkin diagram automorphism)  $\operatorname{Rot}(b_1 \otimes b_2) := \operatorname{Rot}(b_1) \otimes \operatorname{Rot}(b_2), \operatorname{Rot} \circ R = R \circ \operatorname{Rot}$ .
- reversing 01-words (Lusztig involution)  $\vee: B_r \otimes B_s \to B_s \otimes B_r$  with  $b_1 \otimes b_2 \mapsto b_2^{\vee} \otimes b_1^{\vee}$ ,

$$\vee \circ f_i = e_{N-i} \circ \vee, \qquad \vee \circ e_i = f_{N-i} \circ \vee, \qquad \vee \circ R_{r,s} = R_{s,r} \circ \vee.$$

#### Conclusions

### Conclusions

## Additional results

- Eigenvectors of the transfer matrices = idempotents.
- Yang-Baxter algebras, affine nil Temperley-Lieb algebra and Schur polynomials.
- Algorithms to generate vicious/osculating walks.
- The  $\widehat{\mathfrak{sl}}_N$ -Verlinde algebra and cylindric Macdonald functions.

Outlook

- Combinatorial definition of GW invariants and positivity.
- Other Lie algebras.
- Quantum Horn conjecture.
- Categorification of integrable systems.