

# **A Remark on the Hankel Determinant Formula for the Solutions of the Toda Equation**

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# Introduction: Painlevé II Case

## Determinant formula for Yablonski-Vorob'ev polynomials

$$P_{II} : \quad \frac{d^2 u}{dt^2} = 2u^3 - 4tu + 4\left(\alpha + \frac{1}{2}\right)$$

## Yablonski-Vorob'ev polynomials $T_n$ :

$$T_n'' T_n - (T_n')^2 = T_{n+1} T_{n-1} - t T_n^2 : \quad \text{ Toda equation}$$

$$T_0 = 1, \quad T_1 = t,$$

$$T_2 = t^3 - 1,$$

$$T_3 = t^6 - 5t^3 - 5,$$

$$T_4 = t(t^9 - 15t^6 - 175), \dots$$



$$\text{Rational solutions :} \quad u = \frac{d}{dt} \log \frac{T_{n+1}}{T_n}, \quad \alpha = n + \frac{1}{2}$$

# What is Y-V polynomials? → Determinant Formula (K-Ohta 1996)

Jacobi-Trudi Type:

$$T_n \propto \begin{vmatrix} p_N & p_{N+1} & \cdots & p_{2N-1} \\ p_{N-2} & p_{N-1} & \cdots & p_{2N-3} \\ \vdots & \vdots & \ddots & \vdots \\ p_{-N+2} & \cdots & p_0 & p_1 \end{vmatrix}$$

$$\sum_{k=0}^{\infty} p_n(t) \lambda^k = \exp \left[ t\lambda + \frac{\lambda^3}{3} \right]$$

↓

**Specialization of Schur function!**

Hankel Type:

$$T_n = \begin{vmatrix} a_0 & a_1 & \cdots & a_{n-1} \\ a_1 & a_2 & \cdots & a_n \\ \vdots & \vdots & \ddots & \vdots \\ a_{n-1} & a_n & \cdots & a_{2n-2} \end{vmatrix}$$

$$a_0 = t, \quad a_n = a'_{n-1} + \sum_{k=0}^{n-2} a_k a_{n-2-k}$$

↓

**What does this formula mean?**

Generating function for  $a_n$ : log derivative of the Airy fn. (Iwasaki-K-Nakamura, 2002)

$$\frac{\partial}{\partial \lambda} \log \theta(x, t) \sim \sum_{n=0}^{\infty} a_n(x) \lambda^{-n}, \quad \theta(x, \lambda) = \exp \left( -\frac{\lambda^3}{12} \right) \text{Ai} \left( \frac{\lambda^2}{4} - x \right), \quad |\arg \lambda| < \frac{\pi}{2}.$$

# Generic solutions for $P_{II}$ : (K-Masuda 1998, Joshi-K-Mazzocco 2004)

## Hankel determinant formula:

$$\tau_n = \begin{cases} \det(a_{i+j-2})_{i,j \leq n} & n > 0, \\ 1 & n = 0, \\ \det(b_{i+j-2})_{i,j \leq |n|} & n < 0, \end{cases}$$

$$a_n = a'_{n-1} + \psi_{-1} \sum_{k=1}^{n-2} a_k a_{n-2-k}, \quad a_0 = \psi_1,$$

$$b_n = b'_{n-1} + \psi_1 \sum_{k=1}^{n-2} b_k b_{n-2-k}, \quad b_0 = \psi_{-1}$$

$\psi_{\pm 1}$ :

$$\begin{cases} \frac{\psi''_{-1}}{\psi_{-1}} = \frac{\psi''_1}{\psi_1} = -2\psi_{-1}\psi_1 + 2t, \\ \psi'_1\psi_{-1} - \psi_1\psi'_{-1} = 2\alpha. \end{cases}$$

☞  $u_0 = (\log \psi_1)'$  satisfies  $P_{II}[\alpha]$ .

☞  $u_{-1} = -(\log \psi_{-1})'$  satisfies  $P_{II}[\alpha - 1]$ .

☞  $u_N = \left(\log \frac{\tau_{N+1}}{\tau_N}\right)'$  satisfies  $P_{II}[\alpha + N]$ .

## Generating functions: linear problem

$$F(t, \lambda) = \sum_{n=0}^{\infty} a_n (-\lambda)^{-n}, \quad G(t, \lambda) = \sum_{n=0}^{\infty} b_n \lambda^{-n}.$$

↓

## $G$ : Linear problem for $P_{II}[\alpha]$ (Jimbo-Miwa)

$$\boxed{G = \lambda \frac{Y_2}{Y_1}}, \quad \frac{\partial}{\partial \lambda} Y = AY, \quad \frac{\partial}{\partial t} Y = BY, \quad Y = \begin{pmatrix} Y_1 \\ Y_2 \end{pmatrix}$$

$$A = \begin{pmatrix} \frac{1}{4} & 0 \\ 0 & -\frac{1}{4} \end{pmatrix} \lambda^2 + \begin{pmatrix} 0 & -\frac{\psi_1}{2} \\ \frac{\psi_{-1}}{2} & 0 \end{pmatrix} \lambda + \begin{pmatrix} -\frac{z+t}{2} & \frac{\psi_1 u_0}{2} \\ -\frac{\psi_{-1} u_{-1}}{2} & \frac{z+t}{2} \end{pmatrix},$$

$$B = \begin{pmatrix} -\frac{1}{2} & 0 \\ 0 & \frac{1}{2} \end{pmatrix} \lambda + \begin{pmatrix} 0 & \psi_1 \\ -\psi_{-1} & 0 \end{pmatrix}, \quad z = -\psi_{-1}\psi_1$$

## $F$ : Adjoint problem for $P_{II}[\alpha]$

$$\text{ Toda equation: } \quad \frac{dV_n}{dt} = V_n(I_n - I_{n+1}), \quad \frac{dI_n}{dt} = V_{n-1} - V_n$$

**Soliton solution:**  $n \in \mathbb{Z}$

$$V_n = \frac{\tau_{n+1}\tau_{n-1}}{\tau_n^2}, \quad I_n = \frac{d}{dt} \log \frac{\tau_{n-1}}{\tau_n}$$

$$\tau_n = \left( \begin{array}{cccc} f_n^{(1)} & f_{n+1}^{(1)} & \cdots & f_{n+N-1}^{(1)} \\ f_n^{(2)} & f_{n+1}^{(2)} & \cdots & f_{n+N-1}^{(2)} \\ \vdots & \vdots & \cdots & \vdots \\ f_n^{(N)} & f_{n+1}^{(N)} & \cdots & f_{n+N-1}^{(N)} \end{array} \right) \Bigg\} N$$

$$f_n^{(k)} = p_k^n e^{p_k t + p_{k,0}} + p_k^{-n} e^{\frac{1}{p_k} t + p_{k,1}}$$

bilinear equation:

$$\tau_n'' \tau_n - (\tau_n')^2 = \tau_{n+1} \tau_{n-1} - \tau_n^2$$

determinant size = **number of solitons**

**Molecule type solution:**  $0 \leq n (\leq M)$

boundary condition:  $V_0 = 0$  ( $V_M = 0$ )

$$\tau_n = \left( \begin{array}{cccc} a_0 & a_1 & \cdots & a_{N-1} \\ a_1 & a_2 & \cdots & a_N \\ \vdots & \vdots & \cdots & \vdots \\ a_{n-1} & a_n & \cdots & a_{2n-2} \end{array} \right) \Bigg\} n$$

$$\tau_0 = 1, \quad a'_k = a_{k+1} \quad (a_0 = \sum_{j=0}^M c_j e^{\mu_j t})$$

bilinear equation:

$$\tau_n'' \tau_n - (\tau_n')^2 = \tau_{n+1} \tau_{n-1}$$

determinant size = **lattice site**

## Determinant formula of soliton type : (incl. rational solution)

Meaning in the framework of the KP theory: **Clear**

## Determinant formula of molecule type: usually exists for discrete variable

Meaning in the framework of the KP theory: **?**

## Physical meaning of semi-infinite(finite) Toda lattice:

$$\frac{dV_n}{dt} = V_n(I_n - I_{n+1}), \quad \frac{dI_n}{dt} = V_{n-1} - V_n, \quad V_0(= V_N) = 0$$

→ short-circuit the edge of LC ladder circuit →  $V_n(\infty) = I_n(\infty) = 0$

- Not so much interesting physically
- Important in relation to orthogonal functions, Painlevé systems, numerical analysis,...

**Toda equation:**  $\tau_n'' \tau_n - (\tau_n')^2 = \tau_{n+1} \tau_{n-1}, \quad n \in \mathbb{Z}$

Generalization of Hankel determinant formula of molecule type (KMNOY:1998)

$$\frac{\tau_{k+n}}{\tau_k} = \begin{cases} \det(a_{i+j-2}^{(k)})_{i,j=1,\dots,n} & n > 0, \\ 1 & n = 0, \\ \det(b_{i+j-2}^{(k)})_{i,j=1,\dots,|n|} & n < 0, \end{cases}$$

$$\left\{ \begin{array}{l} a_i^{(k)} = a_{i-1}^{(k)'} + \frac{\tau_{k-1}}{\tau_k} \sum_{l=0}^{i-2} a_l^{(k)} a_{i-2-l}^{(k)}, \quad a_0^{(k)} = \frac{\tau_{k+1}}{\tau_k}, \\ b_i^{(k)} = b_{i-1}^{(k)'} + \frac{\tau_{k+1}}{\tau_k} \sum_{l=0}^{i-2} b_l^{(k)} b_{i-2-l}^{(k)}, \quad b_0^{(k)} = \frac{\tau_{k-1}}{\tau_k}, \end{array} \right.$$

...	$\tau_{k-3}$	$\tau_{k-2}$	$\tau_{k-1}$	$\tau_k$	$\tau_{k+1}$	$\tau_{k+2}$	$\tau_{k+3}$	...
...	$\frac{\tau_{k-3}}{\tau_k}$	$\frac{\tau_{k-2}}{\tau_k}$	$\frac{\tau_{k-1}}{\tau_k}$	1	$\frac{\tau_{k+1}}{\tau_k}$	$\frac{\tau_{k+2}}{\tau_k}$	$\frac{\tau_{k+3}}{\tau_k}$	...
...	3×3	2×2	1×1	0×0	1×1	2×2	3×3	...

**$\tau_{k-1} = 0 \longrightarrow$  molecule solution**

# Linear problems and determinant formula

**Auxiliary linear problem:**

$$\Psi'_n = V_{n-1} \Psi_{n-1},$$

$$V_{n-1} \Psi_{n-1} + I_n \Psi_n + \Psi_{n+1} = \lambda \Psi_n,$$

**Adjoint problem:**

$$\Psi_n^{*'} = -V_n \Psi_{n+1}^*,$$

$$\Psi_{n-1}^* + I_n \Psi_n^* + V_{n+1} \Psi_{n+1}^* = \lambda \Psi_n^*$$

**Theorem:**

$$\left[ \frac{\Psi_k}{\Psi_{k+1}} \right]^{(-1)} = \frac{1}{\lambda} \frac{\tau_k}{\tau_{k-1}} \sum_{i=0}^{\infty} b_i^{(k)} \lambda^{-i}$$

$$\left[ \frac{\Psi_{k+1}^*}{\Psi_k^*} \right]^{(-1)} = \frac{1}{\lambda} \frac{\tau_k}{\tau_{k+1}} \sum_{i=0}^{\infty} a_i^{(k)} (-\lambda)^{-i}$$

**Sketch of proof:**

Riccati eq. for  $\Xi_k = \frac{\Psi_k}{\Psi_{k+1}}$ :

$$\frac{\partial \Xi_k}{\partial t} = -V_k \Xi_k^2 + (\lambda - I_k) \Xi_k - 1 \quad (\star)$$

Two kinds of series solutions:

$$\Xi_k^{(-1)} = \frac{1}{\lambda} \sum_{i=0}^{\infty} c_i \lambda^{-i}, \quad \Xi_k^{(1)} = \lambda \sum_{i=0}^{\infty} d_i \lambda^{-i}$$

Recurrence relation for  $b_i$ :

$$b_i^{(k)} = b_{i-1}^{(k)'} + \frac{\tau_{k+1}}{\tau_k} \sum_{l=0}^{i-2} b_l^{(k)} b_{i-2-l}^{(k)}, \quad b_0^{(k)} = \frac{\tau_{k-1}}{\tau_k}$$

$$\rightarrow F = \frac{1}{\lambda} \frac{\tau_k}{\tau_{k-1}} \sum_{i=0}^{\infty} b_i^{(k)} \lambda^{-i} \text{ satisfies } (\star)$$



# Local Lax Pair and Painlevé equations(II,IV & V)

Linear problem for Toda eq.

$$\begin{aligned}\Psi'_{n+1} &= V_n \Psi_n, \\ V_{n-1} \Psi_{n-1} + I_n \Psi_n + \Psi_{n+1} &= \lambda \Psi_n\end{aligned}$$

“Local Lax Pair” (Fadeev-Takhtajan)

$$\tilde{L}_n \phi_n = \phi_{n+1}, \quad \frac{d\phi_n}{dt} = \tilde{B}_n \phi_n,$$

$$\phi_n = \begin{pmatrix} \phi_n^{(1)} \\ \phi_n^{(2)} \end{pmatrix} = e^{-\frac{1}{2}\lambda t} \begin{pmatrix} \Psi_n \\ \frac{\tau_{n-1}}{\tau_{n-2}} \Psi_{n-1} \end{pmatrix}$$

$$\tilde{L}_n(t, \lambda) = \begin{pmatrix} -I_n + \lambda & -\frac{\tau_n}{\tau_{n-1}} \\ \frac{\tau_{n-1}}{\tau_n} & 0 \end{pmatrix},$$

$$\tilde{B}_n(t, \lambda) = \begin{pmatrix} -\frac{1}{2} & 0 \\ 0 & \frac{1}{2} \end{pmatrix} \lambda + \begin{pmatrix} 0 & \frac{\tau_n}{\tau_{n-1}} \\ -\frac{\tau_{n-2}}{\tau_{n-1}} & 0 \end{pmatrix}$$

$P_{II}$ :

$$\frac{d^2 u}{dt^2} = 2u^3 - 4tu + 4 \left( \alpha + \frac{1}{2} \right), \quad u = \frac{d}{dt} \log \frac{\tau_1}{\tau_0}$$

Auxiliary linear problem: ( $\psi_{\pm 1} = \frac{\tau_{\pm 1}}{\tau_0}$ )

$$\frac{\partial Y}{\partial \lambda} = AY, \quad \frac{\partial Y}{\partial t} = BY, \quad Y = \begin{pmatrix} Y_1 \\ Y_2 \end{pmatrix}$$

$$A = \begin{pmatrix} \frac{1}{4} & 0 \\ 0 & -\frac{1}{4} \end{pmatrix} \lambda^2 + \begin{pmatrix} 0 & -\frac{\psi_1}{2} \\ \frac{\psi_{-1}}{2} & 0 \end{pmatrix} \lambda$$

$$+ \begin{pmatrix} -\frac{z+t}{2} & \frac{\psi_1 u_0}{2} \\ -\frac{\psi_{-1} u_{-1}}{2} & \frac{z+t}{2} \end{pmatrix}, \quad z = -\frac{\tau_1 \tau_{-1}}{\tau_0^2}$$

$$B = \begin{pmatrix} -\frac{1}{2} & 0 \\ 0 & \frac{1}{2} \end{pmatrix} \lambda + \begin{pmatrix} 0 & \frac{\tau_1}{\tau_0} \\ -\frac{\tau_{-1}}{\tau_0} & 0 \end{pmatrix},$$

$$\tilde{B}_1(t, \lambda) = B$$

# Description in the framework of the KP theory

$\tau_n(x)$ :  $\tau$  function of 1-dimensional Toda lattice hierarchy and 1st modified KP hierarchy

$(x = (x_1, x_2, \dots), x_1 = t)$

$$D_{x_1} p_{k+1} \left( \frac{1}{2} \tilde{D} \right) \tau_n \cdot \tau_n = p_k \left( \frac{1}{2} \tilde{D} \right) \tau_{n+1} \cdot \tau_{n-1}$$

$$\left[ D_{x_1} p_k \left( \frac{1}{2} \tilde{D} \right) - p_{k+1} \left( \frac{1}{2} \tilde{D} \right) + p_{k+1} \left( -\frac{1}{2} \tilde{D} \right) \right] \tau_{n+1} \cdot \tau_n = 0$$

where  $\sum_{j=0}^{\infty} p_j(x) \lambda^j = \exp \sum_{n=1}^{\infty} x_n \lambda^n$ ,  $\tilde{D} = \left( D_{x_1}, \frac{1}{2} D_{x_2}, \frac{1}{3} D_{x_3}, \dots \right)$ ,  $\tilde{\partial} = \left( \partial_{x_1}, \frac{1}{2} \partial_{x_2}, \frac{1}{3} \partial_{x_3}, \dots \right)$



$$\frac{\tau_{k+n}}{\tau_k} = \begin{cases} \det(a_{i+j-2}^{(k)})_{i,j=1,\dots,n} & n > 0, \\ 1 & n = 0, \\ \det(b_{i+j-2}^{(k)})_{i,j=1,\dots,|n|} & n < 0, \end{cases} \quad \begin{aligned} a_n^{(k)} &= p_n(\tilde{\partial}) \frac{\tau_{k+1}}{\tau_k}, \\ b_n^{(k)} &= (-1)^n p_n(-\tilde{\partial}) \frac{\tau_{k-1}}{\tau_k} \end{aligned}$$

## Sketch of proof:

$$D_{x_1} p_{k+1} \left( \frac{1}{2} \tilde{D} \right) \tau_n \cdot \tau_n = p_k \left( \frac{1}{2} \tilde{D} \right) \tau_{n+1} \cdot \tau_{n-1}$$

$$\left[ D_{x_1} p_k \left( \frac{1}{2} \tilde{D} \right) - p_{k+1} \left( \frac{1}{2} \tilde{D} \right) + p_{k+1} \left( -\frac{1}{2} \tilde{D} \right) \right] \tau_{n+1} \cdot \tau_n = 0$$



Miwa transformation:  $x_j = \frac{l}{j(-\lambda)^j}$



Bäcklund transformation of Toda equation:

$$D_{x_1} \tau_n(l+1) \cdot \tau_n(l) = -\frac{1}{\lambda} \tau_{n+1}(l+1) \tau_{n-1}(l)$$

$$\left( \frac{1}{\lambda} D_{x_1} + 1 \right) \tau_{n+1}(l+1) \cdot \tau_n(l) - \tau_n(l+1) \tau_{n+1}(l) = 0$$



$$\left\{ \begin{array}{l} V_n = \frac{\tau_{n+1}(l) \tau_{n-1}(l)}{\tau_n(l)^2} \\ I_n = \frac{d}{dt} \log \frac{\tau_{n-1}(l)}{\tau_n(l)} \end{array} \right. \quad \left\{ \begin{array}{l} \Psi_{n+1} = \lambda^n \frac{\tau_n(l-1)}{\tau_n(l)} \\ \Psi_n^* = \lambda^{-n} \frac{\tau_n(l+1)}{\tau_n(l)} \end{array} \right.$$



Linear and adjoint linear problems

Now:

$$\frac{\Psi_k}{\Psi_{k+1}} = \frac{1}{\lambda} \frac{\tau_k(l)}{\tau_{k-1}(l)} \sum_{i=0}^{\infty} b_i^{(k)} \lambda^{-i} = \frac{1}{\lambda} \frac{\tau_{k-1}(l-1) \tau_k(l)}{\tau_{k-1}(l) \tau_k(l-1)}$$

$$\rightarrow \sum_{i=0}^{\infty} b_i^{(k)} \lambda^{-i} = \frac{\tau_{k-1}(l-1)}{\tau_k(l-1)} = e^{-\partial_l} \frac{\tau_{k-1}}{\tau_k}$$

$$= \exp \left[ - \sum_{n=0}^{\infty} \frac{(-\lambda)^{-n}}{n} \frac{\partial}{\partial x_n} \right] \frac{\tau_{k-1}}{\tau_k}$$

$$= \sum_{i=0}^{\infty} (-\lambda)^i p_i(-\tilde{\partial}) \frac{\tau_{k-1}}{\tau_k} \rightarrow \boxed{b_i^k = (-1)^i p_i(-\tilde{\partial}) \frac{\tau_{k-1}}{\tau_k}}$$

## Comments:

☞  $\varphi = a_0 = \frac{\tau_{k+1}}{\tau_k}$ ,  $\psi = b_0 = \frac{\tau_{k-1}}{\tau_k}$  satisfy the nonlinear Schrödinger hierarchy:

$$\begin{cases} \varphi_{x_2} = \varphi_{x_1 x_1} + 2\psi\varphi^2, \\ \psi_{x_2} = -\psi_{x_1 x_1} - 2\psi^2\varphi, \end{cases} \quad \begin{cases} \varphi_{x_3} = \varphi_{x_1 x_1 x_1} + 6\psi\varphi\varphi_{x_1}, \\ \psi_{x_2} = \psi_{x_1 x_1 x_1} + 6\psi\varphi\psi_{x_1}, \end{cases} \quad \dots$$

☞ The determinant formula looks like:

$$\frac{\tau_{k+n}}{\tau_k} = \left| \begin{array}{cccc} p_0(\tilde{\partial}) \frac{\tau_{k+1}}{\tau_k} & p_1(\tilde{\partial}) \frac{\tau_{k+1}}{\tau_k} & \cdots & p_{n-1}(\tilde{\partial}) \frac{\tau_{k+1}}{\tau_k} \\ p_1(\tilde{\partial}) \frac{\tau_{k+1}}{\tau_k} & p_2(\tilde{\partial}) \frac{\tau_{k+1}}{\tau_k} & \cdots & p_n(\tilde{\partial}) \frac{\tau_{k+1}}{\tau_k} \\ \vdots & \vdots & \cdots & \vdots \\ p_{n-1}(\tilde{\partial}) \frac{\tau_{k+1}}{\tau_k} & p_n(\tilde{\partial}) \frac{\tau_{k+1}}{\tau_k} & \cdots & p_{2n-2}(\tilde{\partial}) \frac{\tau_{k+1}}{\tau_k} \end{array} \right|_n, \quad p_j(\tilde{\partial}) = \partial_{x_1}^j + \frac{1}{2}\partial_{x_2}\partial_{x_1}^{j-2} + \dots$$

☞ Reduction to 1-dimensional Toda hierarchy is essential (so far) .

☞ The formula resembles the Darboux transformation, but direct relationship is not known, since **the Wronskian structure** is essential to the Darboux transformation.

**Summary: Toda equation:**  $\tau_n'' \tau_n - (\tau_n')^2 = \tau_{n+1} \tau_{n-1}, \quad n \in \mathbb{Z}$

☞ **Hankel determinant formula:**

$$\frac{\tau_{k+n}}{\tau_k} = \begin{cases} \det(a_{i+j-2}^{(k)})_{i,j=1,\dots,n} & n > 0, \\ 1 & n = 0, \\ \det(b_{i+j-2}^{(k)})_{i,j=1,\dots,|n|} & n < 0, \end{cases} \quad \begin{cases} a_i^{(k)} = a_{i-1}^{(k)'} + \frac{\tau_{k-1}}{\tau_k} \sum_{l=0}^{i-2} a_l^{(k)} a_{i-2-l}^{(k)}, & a_0^{(k)} = \frac{\tau_{k+1}}{\tau_k}, \\ b_i^{(k)} = b_{i-1}^{(k)'} + \frac{\tau_{k+1}}{\tau_k} \sum_{l=0}^{i-2} b_l^{(k)} b_{i-2-l}^{(k)}, & b_0^{(k)} = \frac{\tau_{k-1}}{\tau_k}, \end{cases}$$

☞ **Relation to linear problems:**

$\Psi$ : solution of linear problem,  $\Psi^*$ : solution of adjoint problem

$$\frac{\Psi_k}{\Psi_{k+1}} = \frac{1}{\lambda} \frac{\tau_k}{\tau_{k-1}} \sum_{i=0}^{\infty} b_i^{(k)} \lambda^{-i}, \quad \frac{\Psi_{k+1}^*}{\Psi_k^*} = \frac{1}{\lambda} \frac{\tau_k}{\tau_{k+1}} \sum_{i=0}^{\infty} a_i^{(k)} (-\lambda)^{-i}$$

☞ **KP theory:**

$$\frac{\tau_{k+n}}{\tau_k} = \begin{vmatrix} p_0(\tilde{\partial}) \frac{\tau_{k+1}}{\tau_k} & \cdots & p_{n-1}(\tilde{\partial}) \frac{\tau_{k+1}}{\tau_k} \\ \vdots & \cdots & \vdots \\ p_{n-1}(\tilde{\partial}) \frac{\tau_{k+1}}{\tau_k} & \cdots & p_{2n-2}(\tilde{\partial}) \frac{\tau_{k+1}}{\tau_k} \end{vmatrix}$$

☞ **Painlevé equations through the local Lax pair.**

☞ **Finite lattice**