

ISLAND 3 : Algebraic aspects of integrable systems



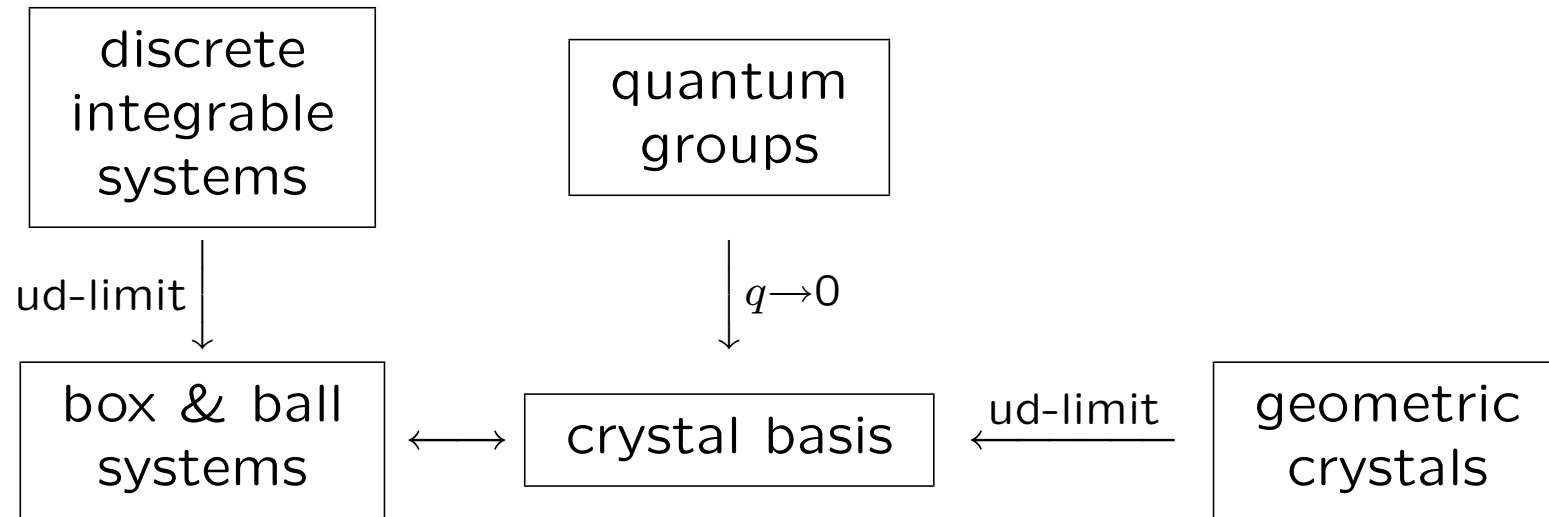
Local Darboux transformations and geometric crystals (II):

Transformations of the dKP equation and their reduction to the dKdV case

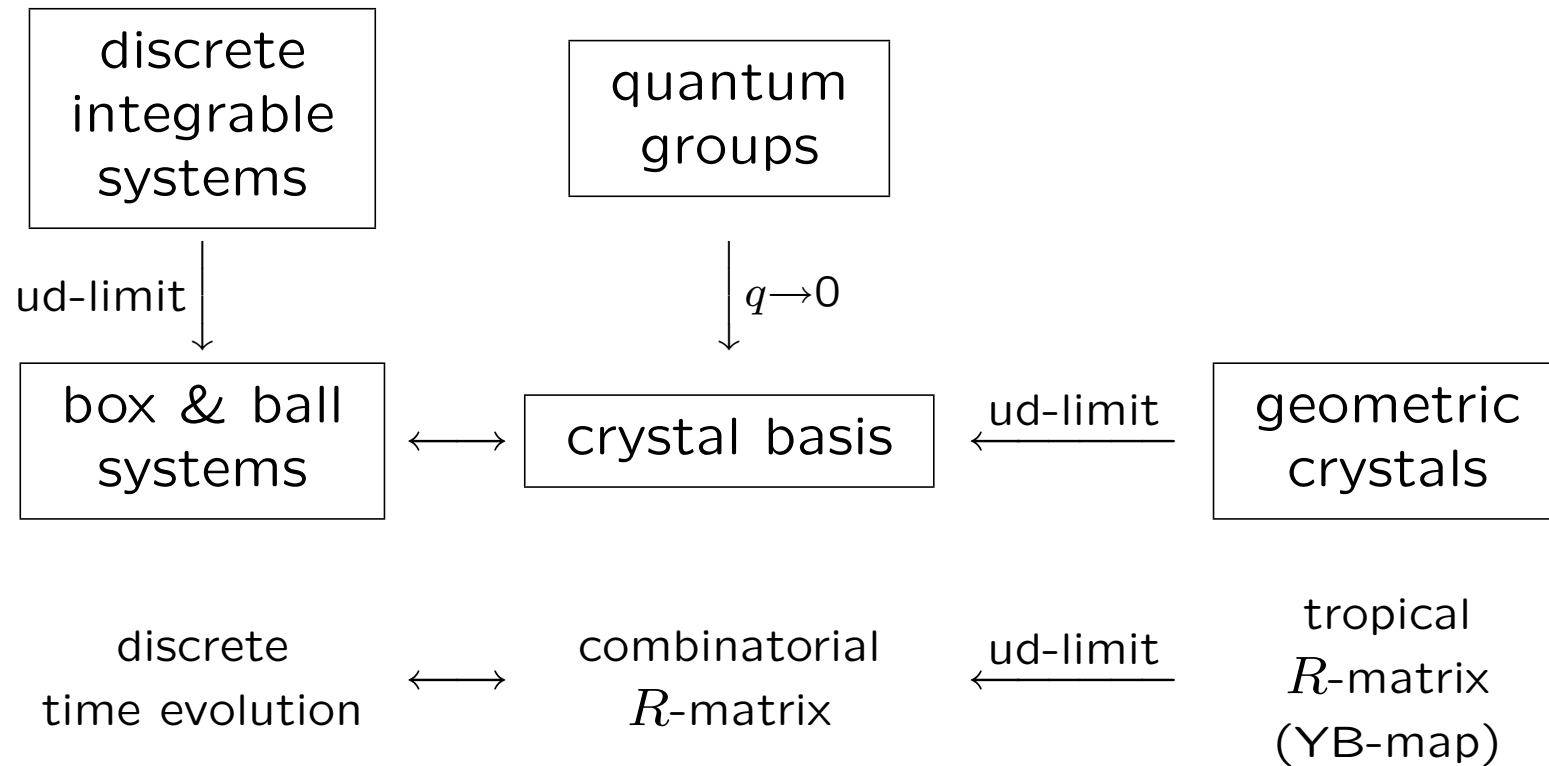
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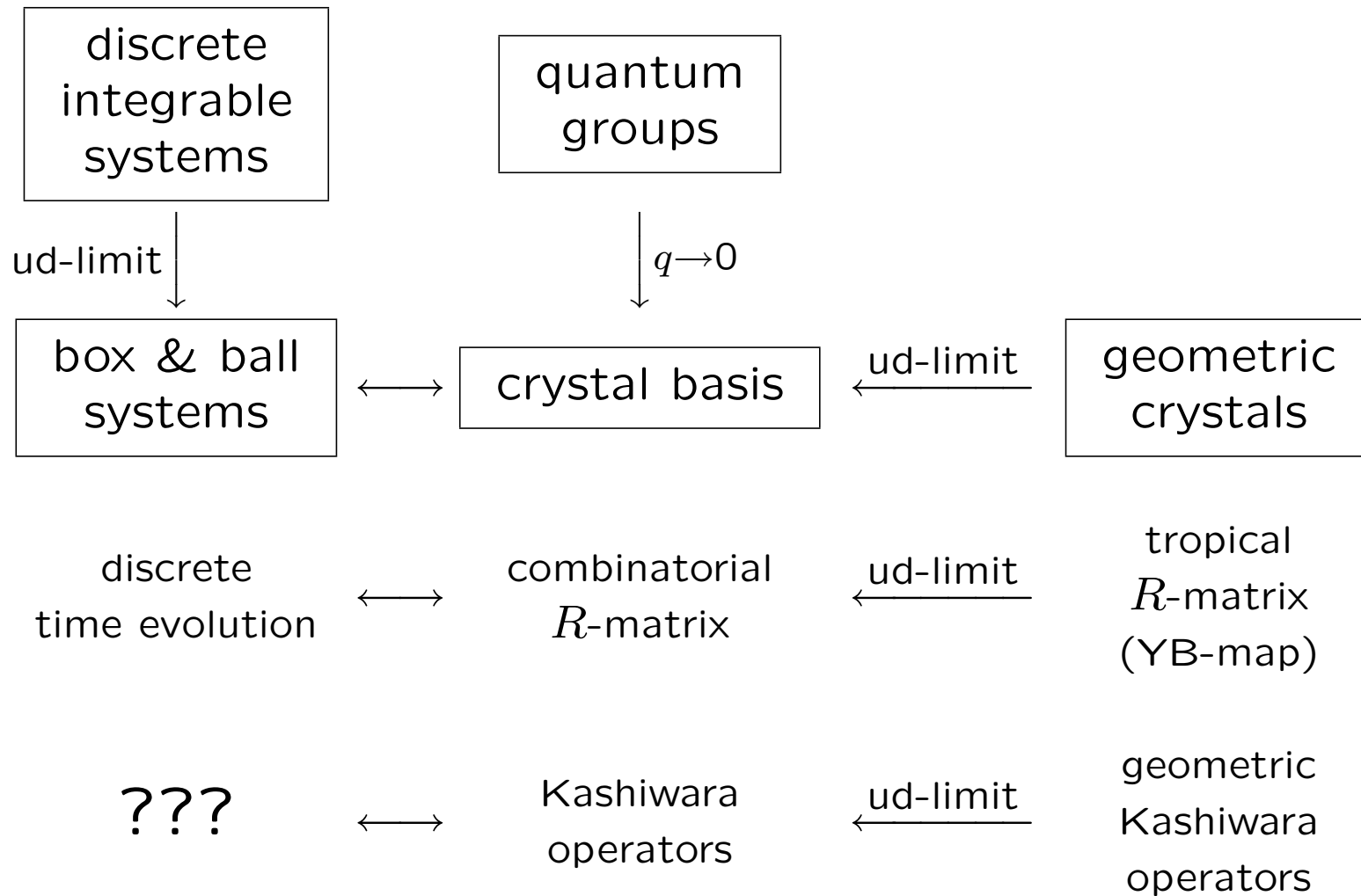
Motivation



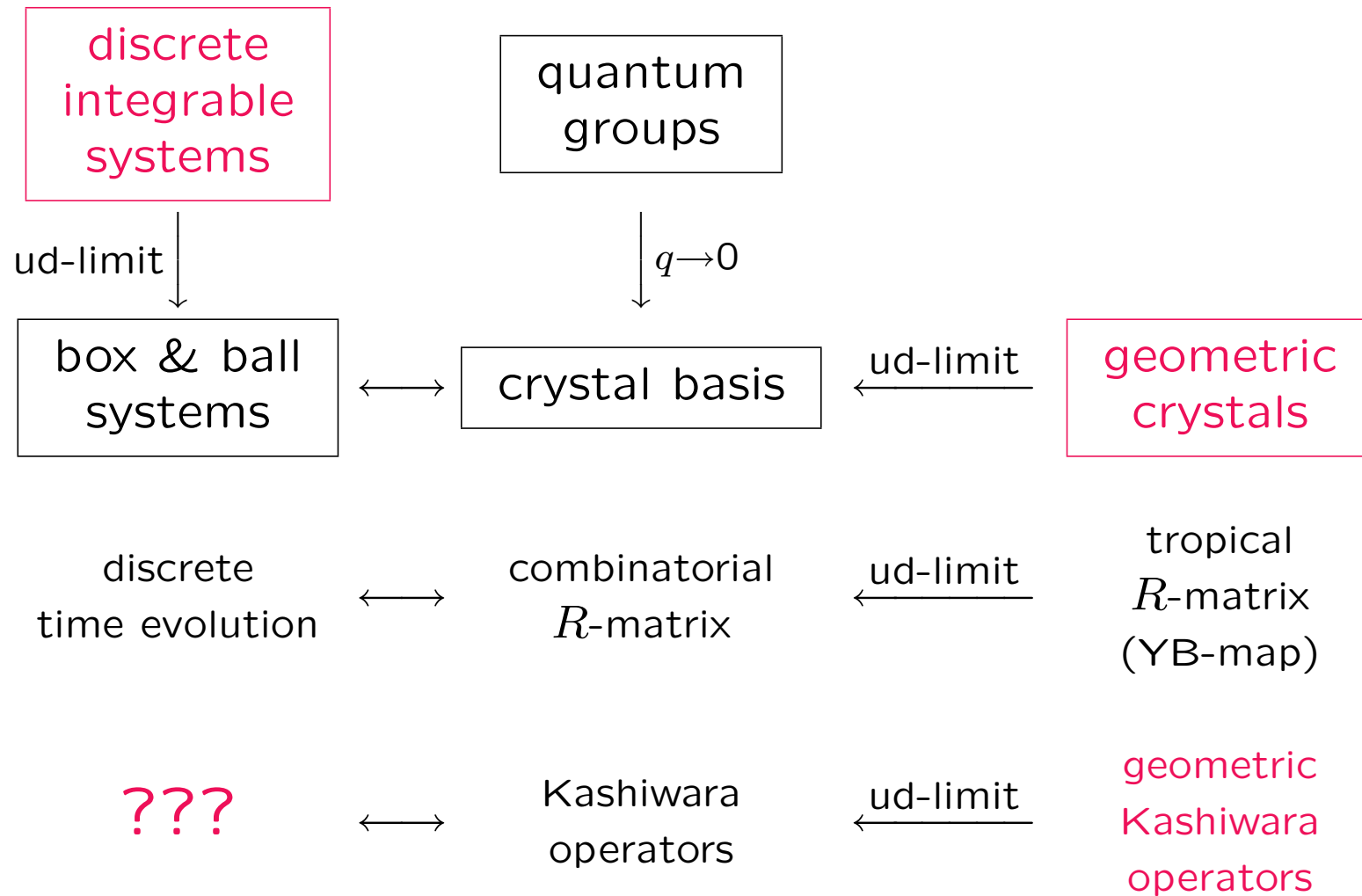
Motivation



Motivation



Motivation



Δ : root system

$\Pi := \{\alpha_i\}_{i \in \mathcal{I}}$: simple roots

$A = (a_{ij})_{i,j \in \mathcal{I}}$, $a_{ij} := \frac{2(\alpha_i, \alpha_j)}{(\alpha_j, \alpha_j)}$: Cartan matrix

X : algebraic variety/ \mathbb{C}

$e_i : \mathbb{C}^* \times X \rightarrow X$ (rational \mathbb{C}^* -actions)

For $c \in \mathbb{C}^*$, $x \in X$, we will use the notation $e_i^c(x) := e_i(c, x)$.

$\gamma_i : X \rightarrow \mathbb{C}$ (rational functions)

$\varepsilon_i : X \rightarrow \mathbb{C}$ (rational functions)

§1. Geometric Crystals (Berenstein-Kazhdan, 2000)

Def. $(X, \{e_i\}_{i \in I}, \{\gamma_i\}_{i \in I}, \{\varepsilon_i\}_{i \in I})$ is a geometric crystal if

$$(i) \quad \gamma_i(e_j^c(x)) = c^{a_{ji}} \gamma_i(x),$$

$$(ii) \quad \varepsilon_i(e_j^c(x)) = c^{-1} \varepsilon_i(x),$$

$$(iii) \quad (0) \quad a_{ij} = a_{ji} = 0 \quad \Rightarrow \quad e_i^{c_1} e_j^{c_2} = e_j^{c_2} e_i^{c_1},$$

$$(1) \quad a_{ij} = a_{ji} = -1 \quad \Rightarrow \quad e_i^{c_1} e_j^{c_1 c_2} e_i^{c_2} = e_j^{c_2} e_i^{c_1 c_2} e_j^{c_1},$$

$$(2) \quad a_{ij} = -2, \quad a_{ji} = -1$$

$$\Rightarrow \quad e_i^{c_1} e_j^{c_1^2 c_2} e_i^{c_1 c_2} e_j^{c_2} = e_j^{c_2} e_i^{c_1 c_2} e_j^{c_1^2 c_2} e_i^{c_1},$$

$$(3) \quad a_{ij} = -3, \quad a_{ji} = -1$$

$$\begin{aligned} \Rightarrow \quad e_i^{c_1} e_j^{c_1^3 c_2} e_i^{c_1^2 c_2} e_j^{c_1^3 c_2^2} e_i^{c_1 c_2} e_j^{c_2} \\ = e_j^{c_2} e_i^{c_1 c_2} e_j^{c_1^3 c_2^2} e_i^{c_1^2 c_2} e_j^{c_1^3 c_2} e_i^{c_1}. \end{aligned}$$

Unipotent Crystal (A_{n-1} -case)

$$GL_n(\mathbb{C}) \begin{cases} \supset \mathcal{B}^+ : \text{upper triangular matrices} \\ \supset \mathcal{U}^- : \text{lower triangular \& diagonal entries} = 1 \end{cases}$$

Notation:

- $E_{ij} \in GL_n, (E_{ij})_{kl} = \delta_{ik}\delta_{jl}$
- $G_j(a) := \sum_i E_{ii} + aE_{j+1,j} = \begin{bmatrix} 1 & & & & & \\ & \dots & & & & \\ & & 1 & & & \\ & & a & 1 & & \\ & & & & \dots & \\ & & & & & 1 \end{bmatrix} \in \mathcal{U}^-$

For $X = \sum_{i \leq j} x_{ij} E_{ij} \in \mathcal{B}^+$ and $c \in \mathbb{C}^*$, we define

- $e_i^c(X) := G_i((c-1)x_{ii}/x_{i,i+1}) X G_i((c^{-1}-1)x_{i+1,i+1}/x_{i,i+1})$,
- $\gamma_i(X) := x_{i+1,i+1}/x_{ii}$

Consider $X \in \mathcal{B}^+$ of the following form:

$$X = \begin{bmatrix} x_1 & 1 & & \\ & x_2 & \cdots & \\ & & \cdots & 1 \\ & & & x_n \end{bmatrix}.$$

We can calculate the action of e_i^c explicitly:

$$e_i^c : \begin{bmatrix} x_1 & 1 & & \\ & x_2 & \cdots & \\ & & \cdots & 1 \\ & & & x_n \end{bmatrix} \mapsto \begin{bmatrix} x_1 & 1 & & \\ & \cdots & \cdots & \\ & & cx_i & \\ & & & 1 \\ & & c^{-1}x_{i+1} & \cdots \\ & & & \cdots & 1 \\ & & & & x_n \end{bmatrix}$$

$$e_i^c : (x_1, \dots, x_i, x_{i+1}, \dots, x_n) \mapsto (x_1, \dots, cx_i, c^{-1}x_{i+1}, \dots, x_n)$$

Consider a product $X \times Y$ ($X, Y \in \mathcal{B}^+$).

We can define an action of e_i^c on $X \times Y$:

$$e_i^c(X \times Y) = e_i^{c_1}(X) \times e_i^{c_2}(Y)$$

where c_1, c_2 are some rational functions w.r.t. x_i, y_i ($i = 1, \dots, n$)

Tropical R -matrix

$$R : \mathcal{B}^+ \times \mathcal{B}^+ \rightarrow \mathcal{B}^+ \times \mathcal{B}^+$$

which satisfies

$$[R, e_i^c] = 0, \quad \forall i$$

Cf. Solvable lattice models \sim Quantum R

$$[R, U_q(\mathfrak{g})] = 0$$

§2. Local Darboux transformation for the discrete KP hierarchy

For $f(\ell, \mathbf{m}) = f(\ell, m_1, \dots, m_M)$:

$$T_\ell f(\ell, \mathbf{m}) = f(\ell + 1, \mathbf{m}),$$

$$T_j f(\ell, \mathbf{m}) = f(\ell, m_1, \dots, m_j + 1, \dots, m_M), \quad j = 1, \dots, M$$

Sato-Wilson operators :

$$W(\ell, \mathbf{m}) := I + w_1 T_\ell^{-1} + w_2 T_\ell^{-2} + \dots,$$

$$\bar{W}(\ell, \mathbf{m}) := \bar{w}_0 + \bar{w}_1 T_\ell^1 + \bar{w}_2 T_\ell^2 + \dots,$$

discrete Sato equations :

$$(T_j \widetilde{W})(1 - \alpha_j + \alpha T_\ell) = B_j \widetilde{W} \quad (\widetilde{W} = W \text{ or } \bar{W}, j = 1, \dots, M),$$

$$B_j = \alpha T_\ell + (1 - \alpha_j) u_j, \quad u_j := \frac{T_j \bar{w}_0}{\bar{w}_0}$$

(Formal) Baker-Akhiezer function :

$$\Psi_\lambda(\ell, \mathbf{m}) := W(\ell, \mathbf{m}) \left(1 - \frac{\lambda}{a}\right)^\ell \prod_{j=1}^M \left(1 - \frac{\lambda}{b_j}\right)^{m_j}$$

that satisfies

$$\begin{aligned} T_j \Psi_\lambda(\ell, \mathbf{m}) &= B_j \Psi_\lambda(\ell, \mathbf{m}) = \left\{ \alpha_j T_\ell + (1 - \alpha_j) u_j \right\} \Psi_\lambda(\ell, \mathbf{m}) \\ \Leftrightarrow \Psi_\lambda &= \frac{1}{a - b_j u_j} \frac{1}{u_j} (a T_\ell - b_j T_j) \Psi_\lambda \end{aligned}$$

[J.J.C. Nimmo, J. Phys. A (1997)]

Compatibility condition : $(T_j B_k) B_j = (T_k B_j) B_k$

$$\begin{cases} u_j(T_j u_k) = u_k(T_k u_j), \\ \alpha_j(1 - \alpha_k)(T_j u_k) + \alpha_k(1 - \alpha_j)(T_\ell u_j) \\ \quad = \alpha_k(1 - \alpha_j)(T_k u_j) + \alpha_j(1 - \alpha_k)(T_\ell u_k) \end{cases}$$

(discrete KP equation)

Linear equations in infinite matrix form

$$T_j \Psi_\lambda(\ell, \mathbf{m}) = \{ \alpha_j T_\ell + (1 - \alpha_j) u_j \} \Psi_\lambda(\ell, \mathbf{m})$$

$$\Leftrightarrow T_j \begin{bmatrix} \vdots \\ \psi(\ell - 1) \\ \psi(\ell) \\ \psi(\ell + 1) \\ \vdots \end{bmatrix} = \begin{bmatrix} \cdots & \cdots & & & & & \\ & (1 - \alpha_j) u_j(\ell - 1) & & & & & \\ & & \alpha_j & & & & \\ & & (1 - \alpha_j) u_j(\ell) & & & & \\ & & & \alpha_j & & & \\ & & & (1 - \alpha_j) u_j(\ell + 1) & \cdots & & \\ & & & & \cdots & & \end{bmatrix} \begin{bmatrix} \vdots \\ \psi(\ell - 1) \\ \psi(\ell) \\ \psi(\ell + 1) \\ \vdots \end{bmatrix}$$

Consider 2×2 subsystem :

$$T_j \begin{bmatrix} \phi(k) \\ \phi(k + 1) \end{bmatrix} = \begin{bmatrix} (1 - \alpha_j) u_j(k) & \alpha_j \\ 0 & (1 - \alpha_j) u_j(k + 1) \end{bmatrix} \begin{bmatrix} \phi(k) \\ \phi(k + 1) \end{bmatrix}.$$

Compatibility of the 2×2 -system follows from that of the infinite case.

Local Darboux transformation

Data :

- Potentials $\{u_j(\ell, \mathbf{m})\}$ that satisfy the dKP equation.
- Eigenfunctions $\{\phi(j)\}_{j \in \mathbb{Z}}$ that satisfy

$$T_j \begin{bmatrix} \phi(k) \\ \phi(k+1) \end{bmatrix} = \begin{bmatrix} (1 - \alpha_j)u_j(k) & \alpha_j \\ 0 & (1 - \alpha_j)u_j(k+1) \end{bmatrix} \begin{bmatrix} \phi(k) \\ \phi(k+1) \end{bmatrix}.$$

Fix an integer k and define $\{\hat{u}_j(\ell, \mathbf{m})\}$ as

$$\hat{u}_j(\ell, \mathbf{m}) = \begin{cases} u_j(\ell, \mathbf{m}) & (\ell \neq k, k+1) \\ u_j(k, \mathbf{m})c_j(k, \mathbf{m}) & (\ell = k) \\ u_j(k+1, \mathbf{m})/c_j(k, \mathbf{m}) & (\ell = k+1) \end{cases}$$

where $c_j(k, \mathbf{m}) := 1 + \phi(k+1) / \{(\alpha_j - 1)\phi(k)u_j(k, \mathbf{m})\}$.

Proposition :

The transformed potentials $\{\hat{u}_j(\ell, \mathbf{m})\}$ satisfy the dKP equation.

§3. Reduction to dKdV

Reduction :
$$\begin{aligned} W(\ell, m_1 + 1, m_2 + 1) &= W(\ell, m_1, m_2), \\ \bar{W}(\ell, m_1 + 1, m_2 + 1) &= \bar{W}(\ell, m_1, m_2) \end{aligned}$$

$$\Rightarrow \Psi_\lambda(\ell, , m_1 + 1, m_2 + 1) = \lambda \Psi_\lambda(\ell, , m_1, m_2)$$

Define $\Phi_\lambda(\ell, , m_1, m_2) := \Psi_\lambda(\ell, , m_1, m_2 + 1)$.

Then $\Phi_\lambda(\ell, , m_1 + 1, m_2) = \lambda \Psi_\lambda(\ell, , m_1, m_2)$.

\Rightarrow 2×2 Lax pair

$$\begin{cases} T_\ell \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \end{bmatrix} = \begin{bmatrix} (1 - \alpha_2^{-1})u & \alpha_2^{-1} \\ \alpha_1^{-1}\lambda & (1 - \alpha_1^{-1})u^{-1} \end{bmatrix} \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \end{bmatrix}, \\ T_1 \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \end{bmatrix} = \begin{bmatrix} (1 - \alpha_1)v & \alpha_1\alpha_2^{-1} \\ \lambda & 0 \end{bmatrix} \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \end{bmatrix}. \end{cases}$$

2×2 Lax pair

$$\begin{cases} T_\ell \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \end{bmatrix} = \begin{bmatrix} (1 - \alpha_2^{-1})u & \alpha_2^{-1} \\ \alpha_1^{-1}\lambda & (1 - \alpha_1^{-1})u^{-1} \end{bmatrix} \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \end{bmatrix}, \\ T_1 \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \end{bmatrix} = \begin{bmatrix} (1 - \alpha_1)v & \alpha_1\alpha_2^{-1} \\ \lambda & 0 \end{bmatrix} \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \end{bmatrix}. \end{cases}$$

Compatibility condition :

$$\begin{cases} (T_1 u)v = u(T_\ell v), \\ v = \frac{1}{T_1 u} - \frac{1 - \alpha_2^{-1}}{1 - \alpha_1^{-1}}u \end{cases}$$

Eliminating v , we obtain the **discrete KdV equation** :

$$\frac{1}{T_1 T_\ell u} - \frac{1}{u} = \delta \{(T_\ell u) - (T_1 u)\}, \quad \delta = \frac{1 - \alpha_2^{-1}}{1 - \alpha_1^{-1}}$$

dKdV as a Box and Ball System

$$\ell \leftrightarrow t, \quad m_1 \leftrightarrow n$$

$$\begin{cases} u_n^{t+1} v_n^t = u_n^t v_{n+1}^t \\ u_n^{t+1} = \frac{1}{v_n^t + \delta u_n^t} \end{cases} \Leftrightarrow \frac{1}{u_{n+1}^{t+1}} - \frac{1}{u_n^t} = \delta \{u_{n+1}^t - u_n^{t+1}\}$$

Ultra-discrete limit [Tokihiko-Takahashi-Matsukidaira-Satsuma, 1996]

$$\lim_{\varepsilon \rightarrow +0} \varepsilon \log \left(e^{A/\varepsilon} + e^{B/\varepsilon} \right) = \max[A, B]$$

$$u_n^t = \exp[U_n^t/\varepsilon], \quad v_n^t = \exp[V_n^t/\varepsilon], \quad \delta = \exp[-1/\varepsilon]$$

$$\begin{cases} e^{(U_n^{t+1} + V_n^t)/\varepsilon} = e^{(U_n^t + V_{n+1}^t)/\varepsilon} \\ e^{U_n^{t+1}/\varepsilon} = \frac{1}{e^{V_n^t/\varepsilon} + e^{(U_n^t - 1)/\varepsilon}} \end{cases} \xrightarrow{\text{ud-lim}} \begin{cases} U_n^{t+1} + V_n^t = U_n^t + V_{n+1}^t \\ U_n^{t+1} = -\max[U_n^t - 1, V_n^t] \end{cases}$$

dKdV as a Yang-Baxter map

$$\text{dKdV} \Leftrightarrow \begin{cases} (T_1 u)v = u(T_\ell v) \\ v = \frac{1}{T_1 u} - \delta u \end{cases} \Leftrightarrow \begin{cases} T_1 u = \frac{1}{v + \delta u} \\ T_\ell v = \frac{v}{u(v + \delta u)} \end{cases}$$

Yang-Baxter map

$$R : (u, v) \mapsto \left(\frac{1}{v + \delta u}, \frac{v}{u(v + \delta u)} \right)$$

This satisfies Yang-Baxter equation :

$$\tilde{R}_{uv} R_{uw} R_{vw} = R_{vw} R_{uw} \tilde{R}_{uv},$$

where $\tilde{R} : (u, v) \mapsto (v, u)$ and

$$R_{vw}(u, v, w) = (u, R(v, w)), \quad \text{etc.}$$

discrete KdV

$$\begin{cases} (T_1 u)v = u(T_\ell v) \\ v = \frac{1}{T_1 u} - \delta u \end{cases} \Leftrightarrow \frac{1}{T_1 T_\ell u} - \frac{1}{u} = \delta \{(T_\ell u) - (T_1 u)\}$$

Cf. [Papageorgiou-Tongas-Veselov, JMP, 2006]

$$\begin{cases} uv = xy \\ u - \frac{\beta_1}{x} = y - \frac{\beta_1}{v} \end{cases} \Leftrightarrow \begin{cases} u = y \frac{\beta_1 + xy}{\beta_2 + xy} \\ v = x \frac{\beta_2 + xy}{\beta_1 + xy} \end{cases}$$

$R : (x, y) \mapsto (u, v)$ satisfies the YB relation.

Local Darboux transformation for dKdV

Proposition :

Assume that the local eigenfunction $\begin{bmatrix} \phi(k) \\ \phi(k+1) \end{bmatrix}$ satisfies

$$T_1 T_2 \begin{bmatrix} \phi(k) \\ \phi(k+1) \end{bmatrix} = \kappa \begin{bmatrix} \phi(k) \\ \phi(k+1) \end{bmatrix}, \quad \kappa \neq 0.$$

Then the local Darboux transformation preserves the reduction condition for the dKdV.

$$\frac{1}{T_1 T_\ell u} - \frac{1}{u} = \delta \{ (T_\ell u) - (T_1 u) \}$$

$$\{u_j(\ell, \mathbf{m})\} \mapsto \hat{u}_j(\ell, \mathbf{m}) = \begin{cases} u_j(\ell, \mathbf{m}) & (\ell \neq k, k+1) \\ u_j(k, \mathbf{m})c_j(k, \mathbf{m}) & (\ell = k) \\ u_j(k+1, \mathbf{m})/c_j(k, \mathbf{m}) & (\ell = k+1) \end{cases}$$

$$c_j(k, \mathbf{m}) := 1 + \phi(k+1) / \{(\alpha_j - 1)\phi(k)u_j(k, \mathbf{m})\}$$

§4. Reduction to dBoussinesq

Reduction 1: $\Psi_\lambda(\ell + 1, m_1 + 1, m_2 + 1) = \lambda \Psi_\lambda(\ell, m_1, m_2)$

Define $\Phi_\lambda(\ell, m_1, m_2, m_3) := \Psi_\lambda(\ell, m_1, m_2 + 1)$,

$\Theta_\lambda(\ell, m_1, m_2, m_3) := \Psi_\lambda(\ell, m_1 + 1, m_2 + 1)$,

Then $\Theta_\lambda(\ell + 1, m_1, m_2) = \lambda \Psi_\lambda(\ell, m_1, m_2)$.

\Rightarrow 3×3 Lax pair

$$\begin{cases} T_\ell \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \\ \Theta_\lambda \end{bmatrix} = \begin{bmatrix} (1 - \alpha_2^{-1})u_2 & \alpha_2^{-1} & 0 \\ 0 & (1 - \alpha_1^{-1})(T_1^{-1}T_\ell^{-1}u_1) & \alpha_1^{-1} \\ \lambda & 0 & 0 \end{bmatrix} \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \\ \Theta_\lambda \end{bmatrix}, \\ T_1 \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \\ \Theta_\lambda \end{bmatrix} = \begin{bmatrix} (1 - \alpha_1)v & \alpha_1\alpha_2^{-1} & 0 \\ 0 & 0 & 1 \\ \alpha_1\lambda & 0 & (1 - \alpha_1)(T_\ell^{-1}u_1) \end{bmatrix} \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \\ \Theta_\lambda \end{bmatrix}. \end{cases}$$

Compatibility condition

$$\begin{cases} u_2(T_\ell v) = (T_1 u_2)v \\ T_\ell v = T_1^{-1}T_\ell^{-1}u_2 - \delta(T_1 u_2) \\ v = u_1 - \delta u_2 \end{cases} \Rightarrow \begin{cases} \frac{T_1 u_2}{u_2} = \frac{T_\ell u_1 - \delta(T_\ell u_2)}{u_1 - \delta u_2} \\ u_1(T_1 u_2) = u_2(T_1^{-1}T_\ell^{-1}u_1) \end{cases}$$

Reduction 2:

$$\Psi_\lambda(\ell, m_1 + 1, m_2 + 1, m_3 + 1) = \lambda \Psi_\lambda(\ell, m_1, m_2, m_3)$$

Define $\Phi_\lambda(\ell, m_1, m_2, m_3) := \Psi_\lambda(\ell, m_1, m_2, m_3 + 1)$,

$$\Theta_\lambda(\ell, m_1, m_2, m_3) := \Psi_\lambda(\ell, m_1, m_2 + 1, m_3 + 1),$$

Then $\Theta_\lambda(\ell, m_1 + 1, m_2, m_3) = \lambda \Psi_\lambda(\ell, m_1, m_2, m_3)$.

\Rightarrow 3×3 Lax pair

$$\left\{ \begin{array}{l} T_\ell \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \\ \Theta_\lambda \end{bmatrix} = \begin{bmatrix} (1 - \alpha_3^{-1})u_3 & \alpha_3^{-1} & \\ & (1 - \alpha_2^{-1})(T_3u_2) & \\ \alpha_1^{-1}\lambda & & (1 - \alpha_1^{-1})(T_2T_3u_1) \end{bmatrix} \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \\ \Theta_\lambda \end{bmatrix}, \\ T_1 \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \\ \Theta_\lambda \end{bmatrix} = \begin{bmatrix} (1 - \alpha_1)v_3 & \alpha_1\alpha_3^{-1} & \\ & (1 - \alpha_1)v_2 & \alpha_1\alpha_2^{-1} \\ \lambda & & \end{bmatrix} \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \\ \Theta_\lambda \end{bmatrix}. \end{array} \right.$$

3 × 3 Lax pair

$$\left\{ \begin{array}{l} T_\ell \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \\ \Theta_\lambda \end{bmatrix} = \begin{bmatrix} (1 - \alpha_3^{-1})u_3 & \alpha_3^{-1} & \\ & (1 - \alpha_2^{-1})/(u_1u_3) & \alpha_2^{-1} \\ \alpha_1^{-1}\lambda & & (1 - \alpha_1^{-1})u_1 \end{bmatrix} \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \\ \Theta_\lambda \end{bmatrix}, \\ T_1 \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \\ \Theta_\lambda \end{bmatrix} = \begin{bmatrix} (1 - \alpha_1)v_3 & \alpha_1\alpha_3^{-1} & \\ & (1 - \alpha_1)v_2 & \alpha_1\alpha_2^{-1} \\ \lambda & & \end{bmatrix} \begin{bmatrix} \Psi_\lambda \\ \Phi_\lambda \\ \Theta_\lambda \end{bmatrix}. \end{array} \right.$$

Compatibility condition :

$$\left\{ \begin{array}{l} (T_1u_3)v_3 = u_3(T_\ell v_3), \quad u_1u_3v_2 = (T_1u_1)(T_1u_3)(T_\ell v_2), \\ v_3 = (T_1u_1) - \delta_3u_3, \quad v_2 = \frac{(T_1u_1)(T_1u_3)}{u_3} - \frac{\delta_2}{u_1u_3}. \end{array} \right.$$

(modified d-Boussinesq equation (?))

m-dBoussinesq as Yang-Baxter map

(Cf. Papageorgiou-Tongas-Veselov, JMP, 2006)

$$\text{Compatibility} \Leftrightarrow \begin{cases} T_1 u_1 = v_2 + \delta_2 u_2, & T_1 u_2 = \frac{u_1 u_2 v_1 + \delta_1}{u_1 (v_2 + \delta_2 u_2)}, \\ T_\ell v_1 = \frac{u_1 u_2 v_1 v_2 + \delta_1 v_2}{u_1 u_2 (v_2 + \delta_2 u_2)}, & T_\ell v_2 = \frac{u_1^2 u_2 v_1}{u_1 u_2 v_1 + \delta_1} \end{cases}$$

Yang-Baxter map

$$R : (u_1, u_2, v_1, v_2)$$

$$\mapsto \left(v_2 + \delta_2 u_2, \frac{u_1 u_2 v_1 + \delta_1}{u_1 (v_2 + \delta_2 u_2)}, \frac{u_1 u_2 v_1 v_2 + \delta_1 v_2}{u_1 u_2 (v_2 + \delta_2 u_2)}, \frac{u_1^2 u_2 v_1}{u_1 u_2 v_1 + \delta_1} \right)$$

This satisfies Yang-Baxter equation :

$$\tilde{R}_{uv} R_{uw} R_{vw} = R_{vw} R_{uw} \tilde{R}_{uv},$$

where $\tilde{R} : (u_1, u_2, v_1, v_2) \mapsto (v_1, v_2, u_1, u_2)$ and

$$R_{vw}(u, v, w) = (u, R(v, w)), \quad \text{etc.}$$

Local Darboux transformation for m-dBoussinesq

Proposition :

Assume that the local eigenfunction $\begin{bmatrix} \phi(k) \\ \phi(k+1) \end{bmatrix}$ satisfies

$$T_1 T_2 T_3 \begin{bmatrix} \phi(k) \\ \phi(k+1) \end{bmatrix} = \kappa \begin{bmatrix} \phi(k) \\ \phi(k+1) \end{bmatrix}, \quad \kappa \neq 0.$$

Then the local Darboux transformation preserves the reduction condition for the dBoussinesq.

$$\{u_j(\ell, \mathbf{m})\} \mapsto \hat{u}_j(\ell, \mathbf{m}) = \begin{cases} u_j(\ell, \mathbf{m}) & (\ell \neq k, k+1) \\ u_j(k, \mathbf{m})c_j(k, \mathbf{m}) & (\ell = k) \\ u_j(k+1, \mathbf{m})/c_j(k, \mathbf{m}) & (\ell = k+1) \end{cases}$$

$$c_j(k, \mathbf{m}) := 1 + \phi(k+1) / \{(\alpha_j - 1)\phi(k)u_j(k, \mathbf{m})\}$$