Linear deformations of the matrix product and integrable ODEs.

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References

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3. Odesskii A.V. and Sokolov V. V. Integrable matrix equations related to pairs of compatible associative algebras, Journal Phys. A, 2006, **39**, 12447–12456.

4. Odesskii A. V. and Sokolov V. V. *Pairs of compatible associative algebras, classical Yang-Baxter equation and quiver representations*, accepted to Comm. in Math. Phys. 2007, , no., –. Consider the following class of matrix ODEs

$$\frac{dx}{dt} = [x, Q(x)],$$

where x is $n \times n$ matrix, Q is a constant linear operator $Q: Mat_n \rightarrow Mat_n$.

For instance, the matrix equation

$$\frac{dx}{dt} = [x, xc + cx] = x^2 c - c x^2.$$
(1)

is integrable for any constant matrix c and any n. This equation possesses an infinite set of homogeneous integrals $H_{i,j}$, where i and j are degrees with respect to x and c.

For example,

$$H_{1,1} = \text{trace}(xc), \quad H_{2,1} = \text{trace}(x^2c).$$

Equation (1) is Hamiltonian one with respect to the standard matrix linear Poisson bracket, given by the Hamiltonian operator ad_x , and Hamiltonian function $H_{2,1}$.

Matrix equations of arbitrary size like (1) are important because of a possibility to make different reductions.

For the most trivial reduction one may regard x as a block-matrix. In this case (1) becomes a system of several matrix equations for the block entries of x.

Under reduction $x^T = -x, c^T = c$ equation (1) is a commuting flow for the *n*-dimensional Euler equation.

Let x and c in the equation $\frac{dx}{dt} = x^2 \, c - c \, x^2$ be represented by matrices of the form

$$x = \begin{pmatrix} 0 & u_1 & 0 & 0 & \cdot & 0 \\ 0 & 0 & u_2 & 0 & \cdot & 0 \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ 0 & 0 & 0 & 0 & \cdot & u_{N-1} \\ u_N & 0 & 0 & 0 & \cdot & 0 \end{pmatrix}, \ c = \begin{pmatrix} 0 & 0 & 0 & \cdot & 0 & J_N \\ J_1 & 0 & 0 & \cdot & 0 & 0 \\ 0 & J_2 & 0 & \cdot & 0 & 0 \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ 0 & 0 & 0 & \cdot & J_{N-1} & 0 \end{pmatrix}$$

,

where u_k and J_k are block matrices (of any dimension). It follows from equation (1) that u_k satisfy the nonabelian Volterra equation

$$\frac{d}{dt}u_k = u_k \circ u_{k+1} \circ J_{k+1} - J_{k-1} \circ u_{k-1} \circ u_k.$$

A multiplication \circ defined on the vector space Mat_n of all $n \times n$ matrices is said to be compatible with the matrix multiplication if the product

$$X \bullet Y = XY + \lambda X \circ Y \tag{2}$$

is associative for any constant λ .

Since the matrix algebra is rigid, the multiplication (2) is isomorphic to the matrix multiplication for almost all values of the parameter λ . This means that there exists a formal series of the form

$$S_{\lambda} = 1 + R \ \lambda + O(\lambda^2), \tag{3}$$

with the coefficients being linear operators on Mat_n , such that

$$S_{\lambda}(X) S_{\lambda}(Y) = S_{\lambda}(XY + \lambda \ X \circ Y).$$
(4)

It follows from this formula that the multiplication \circ is given by

$$X \circ Y = R(X)Y + XR(Y) - R(XY).$$
 (5)

where $R: Mat_n \rightarrow Mat_n$ is a linear operator.

Example 1. Let c be an element of Mat_n and R: $X \to cX$ be the operator of left multiplication by c. Then the corresponding multiplication $X \circ Y = X cY$ is associative and compatible with matrix multiplication.

Example 2. Suppose $a, b \in Mat_2$; then the product

$$X \circ Y = (aX - Xa)(bY - Yb)$$

is compatible with the standard product in Mat_2 . The corresponding operator R is given by R(X) = aXb - abX.

Classical associative Yang-Baxter equation

Let $\lambda \to S_{\lambda}$ be a meromorphic function with values in $End(Mat_n)$ such that $S_0 = Id$ and

$$S_{\lambda}(X) S_{\lambda}(Y) = S_{\lambda}(XY + \lambda X \circ Y),$$

where \circ is an associative multiplication compatible with the matrix one.

Theorem. The formula

$$r(u,v) = \frac{1}{u-v} S_u S_v^{-1}$$

defines a solution to the associative Yang-Baxter equation

$$(r(u,w)x)(r(u,v)y) - r(u,v)((r(v,w)x)y) =$$
$$r(u,w)(x(r(w,v)y)).$$

Integrable matrix ODEs.

It is convenient to write down the operator \boldsymbol{R} in the form

$$R(x) = a_1 x b^1 + \dots + a_p x b^p + c x,$$
 (6)

where $a_i, b^i, c \in Mat_n$, with p being smallest possible.

Let \circ be a multiplication (5) compatible with the matrix product. Consider the following matrix differential equation

$$\frac{dx}{dt} = [R(x) + R^*(x), x],$$
(7)

where

$$R^*(x) = b^1 x a_1 + \dots + b^p x a_p + x c.$$

It turns out that equation (7) possesses the Lax representation

$$\frac{dL}{dt} = [A, L],$$

where

$$L = \left(S_{\lambda}^{-1}\right)^*(x), \qquad A = \frac{1}{\lambda}S_{\lambda}(x). \tag{8}$$

As usual, the integrals of motion for (7) are given by coefficients of different powers of λ in trace (L^k) , k = 1, 2...

To make these formulas constructive, we should find S_{λ} and S_{λ}^{-1} in a closed form i.e. as analytic operator-valued functions.

In the case of Example 1, we have

$$R(x) = cx, \qquad S_{\lambda}(x) = (1 + \lambda c) x,$$
$$L = x(1 + \lambda c)^{-1}.$$

The Lax equation is equivalent to the well known integrable matrix equation

$$\frac{dx}{dt} = [x, xc + cx] = x^2 c - c x^2.$$

The *L*-operator produces an infinite set of homogeneous integrals $H_{i,j}$, where *i* and *j* are degrees with respect to *x* and *c*. The simplest are

$$H_{1,1} = \text{trace}(xc), \quad H_{2,1} = \text{trace}(x^2c).$$

A generalization of Example 2 is defined by two arbitrary constant matrices A and B such that

$$A^2 = B^2 = 1. (9)$$

The corresponding associative multiplication is given by (5), where

$$R(x) = AxB + BAx.$$
(10)

This leads to the following integrable matrix equation

$$x_t = [x, BxA + AxB + xBA + BAx].$$
(11)

The Lax representation (8) for (11) is given by the following explicit formulas for S_{λ} and S_{λ}^{-1} :

$$S_{\lambda}(x) = \frac{1-q}{2}BxB + \frac{1+q}{2}x + \lambda(AxB + BAx),$$
$$S_{\lambda}^{-1}(x) = \frac{1}{q}(1+\lambda K)^{-1}\left(\frac{q-1}{2}BxB + \frac{1+q}{2}x + \lambda(ABx - AxB)\right),$$
where $q = \sqrt{1-4\lambda^2}, \ K = AB + BA.$

The simplest linear and quadratic first integrals for (11) generated by the *L*-operator are given by

 $H_{1,1} = \operatorname{trace} [x(AB+BA)], \qquad H_{2,1} = \operatorname{trace} [ABx^2 + AxBx].$

Equation (11) admits the following skew- symmetric reduction

$$x^T = -x, \qquad B = A^T. \tag{12}$$

Different integrable so(n)-models provided by reduction (12) are in one-to-one correspondence with equivalence classes with respect to the SO(n) gauge action of $n \times n$ matrices A such that $A^2 = 1$.

For the real matrix A, a canonical form for such equivalence class can be chosen as

$$A = \begin{pmatrix} \mathbf{1}_p & T \\ \mathbf{0} & \mathbf{1}_{n-p} \end{pmatrix}.$$
 (13)

Here $\mathbf{1}_s$ stands for the unity $s \times s$ matrix and $T = \{t_{ij}\}$, where $t_{ij} = \delta_{ij}\alpha_i$. This canonical form is defined by the discrete natural parameter p and continuous parameters $\alpha_1, \ldots, \alpha_r$, where $r = \min(p, n - p)$.

For example, in the case n = 4 the equivalence classes with p = 2 and p = 1 give rise to the Steklov and the Poincare integrable models, correspondingly.

Thus, whereas (1) is a matrix version of the so(4)Schottky-Manakov top, equation (11)-(13) with p = [n/2] and p = 1 can be regarded as new so(n) generalizations for the so(4) Steklov and Poincare models. A generalization of the basic matrix integrable model (7):

$$\frac{dx}{dt} = [x, v] + x \star x,$$
$$\frac{dv}{dt} = [x, u] + x \star v,$$

$$\frac{du}{dt} = x \star u,$$

where $x, v, w \in Mat_n$ and

$$X \star Y = [R(X), Y] + [R^T(Y), X] + R^T([X, Y]).$$

If $v = u = 0$, we get

$$\frac{dx}{dt} = [R(x) + R^*(x), x].$$

Algebraic structures.

Let us write R in the form

$$R(x) = a_1 X b^1 + \dots + a_p X b^p + c X$$
(14)

with minimal p. Here $a_i, b^i, c \in Mat_N$.

Lemma 1. If the multiplication defined by

$$X \circ Y = R(X)Y + XR(Y) - R(XY).$$

is associative, then

$$a_i a_j = \phi_{i,j}^k a_k + \mu_{i,j} \mathbf{1}, \qquad b^i b^j = \psi_k^{i,j} b^k + \lambda^{i,j} \mathbf{1}$$

for some tensors $\phi_{i,j}^k, \mu_{i,j}, \psi_k^{i,j}, \lambda^{i,j}$.

This means that the vector spaces spanned by $1, a_1, \ldots a_p$ and $1, b_1, \ldots b_p$ are associative algebras. We denote them by \mathcal{A} and \mathcal{B} . The algebras \mathcal{A} and \mathcal{B} have to be related by certain consistency conditions.

The simplest example of a similar structure can be described as follows.

Associative bi-algebras.

Let \mathcal{A} and \mathcal{B} be associative algebras with basis A_1, \ldots, A_p and B^1, \ldots, B^p and structural constants $\phi_{j,k}^i$ and $\psi_{\gamma}^{\alpha,\beta}$. Suppose that the structural constants satisfy the following identities:

$$\phi_{j,k}^{s}\psi_{s}^{l,i} = \phi_{s,k}^{l}\psi_{j}^{s,i} + \phi_{j,s}^{i}\psi_{k}^{l,s}, \qquad 1 \le i, j, k, l \le p.$$

Then the algebra \mathcal{M} of dimension $2p + p^2$ with the basis $A_i, B^j, A_i B^j$ and relations

$$B^i A_j = \psi_j^{k,i} A_k + \phi_{j,k}^i B^k$$

is associative. Note that $A_i B^j$ form an associative subalgebra of dimension p^2 . Consider the vector space \mathcal{L} spanned by A_i, B^j . It is clear that \mathcal{A} and \mathcal{B} act on \mathcal{L} by right and left multiplication, correspondingly.

Let (\cdot, \cdot) be a non-degenerate symmetric scalar product on the space \mathcal{L} such that

$$(A_i, A_j) = (B^i, B^j) = 0, \qquad (A_i, B^j) = \delta_i^j.$$

Then the consistency condition means that

 $(b_1b_2, v) = (b_1, b_2v), \quad (v, a_1a_2) = (va_1, a_2)$ for any $a_1, a_2 \in \mathcal{A}, \ b_1, b_2 \in \mathcal{B}$ and $v \in \mathcal{L}$.

General algebraic structure.

Definition. By week \mathcal{M} -structure on a linear space \mathcal{L} we mean a collection of the following data:

- Two subspaces \mathcal{A} and \mathcal{B} and distinguished element $1 \in \mathcal{A} \cap \mathcal{B} \subset \mathcal{L}$.
- Associative products $\mathcal{A} \times \mathcal{A} \to \mathcal{A}$ and $\mathcal{B} \times \mathcal{B} \to \mathcal{B}$ with unity 1.
- Left action $\mathcal{B} \times \mathcal{L} \to \mathcal{L}$ of the algebra \mathcal{B} and right action $\mathcal{L} \times \mathcal{A} \to \mathcal{L}$ of the algebra \mathcal{A} on the space \mathcal{L} , which commute to each other.
- A non-degenerate symmetric scalar product (\cdot, \cdot) on the space \mathcal{L} .

These data should satisfy the following properties:

1. dim $\mathcal{A} \cap \mathcal{B} = \dim \mathcal{L}/(\mathcal{A} + \mathcal{B}) = 1$. Intersection of \mathcal{A} and \mathcal{B} is a one dimensional space spanned by the unity 1.

2. Restriction of the action $\mathcal{B} \times \mathcal{L} \to \mathcal{L}$ to subspace $\mathcal{B} \subset \mathcal{L}$ is the product in \mathcal{B} . Restriction of the action $\mathcal{L} \times \mathcal{A} \to \mathcal{L}$ to subspace $\mathcal{A} \subset \mathcal{L}$ is the product in \mathcal{A} .

3.
$$(a_1, a_2) = (b_1, b_2) = 0$$
 and
 $(b_1b_2, v) = (b_1, b_2v), \quad (v, a_1a_2) = (va_1, a_2)$
for any $a_1, a_2 \in \mathcal{A}, \ b_1, b_2 \in \mathcal{B}$ and $v \in \mathcal{L}$.

It follows from these properties that (\cdot, \cdot) gives a nondegenerate pairing between $\mathcal{A}/\mathbb{C}1$ and $\mathcal{B}/\mathbb{C}1$, so dim $\mathcal{A} =$ dim \mathcal{B} and dim $\mathcal{L} = 2 \dim \mathcal{A}$. For given week \mathcal{M} -structure we can define an universal associative algebra generated by \mathcal{L} .

Definition. By week \mathcal{M} -algebra associated with week \mathcal{M} -structure on \mathcal{L} we mean an associative algebra $U(\mathcal{L})$ with the following properties:

1. $\mathcal{L} \subset \mathcal{U}(\mathcal{L})$ and the actions $\mathcal{B} \times \mathcal{L} \to \mathcal{L}$, $\mathcal{L} \times \mathcal{A} \to \mathcal{L}$ are restrictions of the product in $U(\mathcal{L})$.

2. For any algebra X with the property **1** there exists a unique homomorphism of algebras $X \to U(\mathcal{L})$, which is identical on \mathcal{L} .

It is easy to see that if $U(\mathcal{L})$ exists, then it is unique for given \mathcal{L} .

Let us describe the structure of $U(\mathcal{L})$ explicitly.

Let $\{1, A_1, ..., A_p\}$ be a basis of \mathcal{A} and $\{1, B^1, ..., B^p\}$ be a dual basis of \mathcal{B} (which means that $(A_i, B^j) = \delta_i^j$). Let $C \in \mathcal{L}$ doesn't belong to the sum of \mathcal{A} and \mathcal{B} . Since (\cdot, \cdot) is non- degenerate, we have $(1, C) \neq 0$. Multiplying Cby constant, we may assume that (1, C) = 1. Adding linear combination of $1, A_1, ..., A_p, B^1, ..., B^p$ to C, we can assume that $(C, C) = (C, A_i) = (C, B^j) = 0$. Such element C is uniquely determined by choosing basis in \mathcal{A} and \mathcal{B} . Let us define an element $K \in U(\mathcal{L})$ by the formula $K = A_i B^i + C$.

Definition. A week \mathcal{M} -structure on \mathcal{L} is called \mathcal{M} structure if $K \in U(\mathcal{L})$ is a central element of the algebra $U(\mathcal{L})$.

Theorem 1. The algebra $U(\mathcal{L})$ is spanned by the elements K^s , A_iK^s , B_jK^s , $A_iB^jK^s$, where i, j = 1, ..., p, and s = 0, 1, 2, ...

Theorem 2. For any representation $U(\mathcal{L}) \rightarrow Mat_n$ given by

$$A_1 \to a_1, ..., A_p \to a_p, B^1 \to b^1, ..., B^p \to b^p, C \to c$$

formula (5) defines associative product on Mat_n compatible with the usual product.

Example 3. Suppose \mathcal{A} and \mathcal{B} are generated by elements $a \in \mathcal{A}$ and $b \in \mathcal{B}$ such that $a^{p+1} = b^{p+1} = 1$. Assume that $(b^i, a^{-i}) = \epsilon^i - 1$, (1, c) = 1 and other scalar products are equal to zero. Here ϵ is a primitive root of unity of order p. Let

$$b^{i}a^{j} = \frac{\epsilon^{-j} - 1}{\epsilon^{-i-j} - 1}a^{i+j} + \frac{\epsilon^{i} - 1}{\epsilon^{i+j} - 1}b^{i+j}$$

for $i + j \neq 0$ modulo p and

$$b^{i}a^{-i} = 1 + (\epsilon^{i} - 1)c, \quad ca^{i} = \frac{1}{1 - \epsilon^{i}}a^{i} + \frac{1}{\epsilon^{i} - 1}b^{i},$$
$$b^{i}c = \frac{1}{\epsilon^{-i} - 1}a^{i} + \frac{1}{1 - \epsilon^{-i}}b^{i}$$

for $i \neq 0$ modulo p. These formulas define an \mathcal{M} -structure.

The central element has the following form

$$K = c + \sum_{0 < i < p} \frac{1}{\epsilon^i - 1} a^{-i} b^i.$$

Let *a*, *t* be linear operators in some vector space. Assume that $a^{p+1} = 1$, $at = \epsilon ta$ and the operator t - 1 is invertible. It is easy to check that the formulas

$$A \to a, \qquad B \to \frac{\epsilon t - 1}{t - 1}a, \qquad C \to \frac{t}{t - 1}$$

define a representation of the algebra $U(\mathcal{L})$.

Case of semi-simple algebras \mathcal{A} and \mathcal{B}

Suppose a vector space ${\mathcal L}$ is equipped with a weak ${\mathcal M}\text{-}$ structure such that

 $\mathcal{A} = \bigoplus_{1 \leq i \leq r} End(V_i), \qquad \mathcal{B} = \bigoplus_{1 \leq j \leq s} End(W_j),$ $\dim V_i = m_i, \qquad \dim W_j = n_j.$ Then \mathcal{L} as $\mathcal{A} \otimes \mathcal{B}$ -module is isomorphic to $\bigoplus_{1 \leq i \leq r, 1 \leq j \leq s} (V_i^* \otimes W_j)^{a_{i,j}}$, for some $a_{i,j} \geq 0$

Theorem. For any i, j

$$\sum_{j=1}^{s} a_{i,j} n_j = 2m_i, \qquad \sum_{i=1}^{r} a_{i,j} m_i = 2n_j.$$
(15)

The matrix of linear system (15) has the form

$$Q = \begin{pmatrix} 2 & -A \\ -A^t & 2 \end{pmatrix}.$$

According to the result by E. Vinberg, if the kernel of indecomposable matrix Q contains an integer positive vector, them Q is the Cartan matrix of an affine Dynkin diagram.

Moreover, it follows from the structure of Q that this diagram is a simple laced affine Dynkin diagram with a partition of the set of vertices into two subsets such that vertices of the same subset are not connected.

Theorem. Let A be an $r \times s$ matrix of multiplicities for an indecomposable weak \mathcal{M} -structure. Then, after a possible permutation of rows and columns and the transposition, a matrix A is equal to one in the following list:

1. A = (2). Here r = s = 1, $n_1 = m_1 = m$. The corresponding Dynkin diagram is of the type \tilde{A}_1 .

2. $a_{i,i} = a_{i,i+1} = 1$ and $a_{i,j} = 0$ for other pairs i, j. Here $r = s = k \ge 2$, the indexes are taken modulo k, and $n_i = m_i = m$. The corresponding Dynkin diagram is \tilde{A}_{2k-1} .





3. $A = \begin{pmatrix} 1 & 1 & 0 & 0 \\ 1 & 0 & 1 & 0 \\ 1 & 0 & 0 & 1 \end{pmatrix}$. Here r = 3, s = 4 and $n_1 = 3m$, $n_2 = n_3 = n_4 = m$, $m_1 = m_2 = m_3 = 2m$. The Dynkin diagram is \tilde{E}_6 :



4.
$$A = \begin{pmatrix} 1 & 1 & 0 & 0 & 0 \\ 0 & 1 & 1 & 1 & 0 \\ 0 & 0 & 0 & 1 & 1 \end{pmatrix}.$$

Here r = 3, s = 5 and

$$n_1 = m, \quad n_2 = 3m, \quad n_3 = 2m, \quad n_4 = 3m, n_5 = m, \quad m_1 = 2m, \quad m_2 = 4m, \quad m_3 = 2m.$$

The Dynkin diagram is \tilde{E}_7 :



5.
$$A = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 \\ 1 & 1 & 1 & 0 & 0 \\ 0 & 0 & 1 & 1 & 0 \\ 0 & 0 & 0 & 1 & 1 \end{pmatrix}$$

Here $n_1 = 4m$, $n_2 = 3m$, $n_3 = 5m$, $n_4 = 3m$, $n_5 = m$, $m_1 = 2m$, $m_2 = 6m$, $m_3 = 4m$, $m_4 = 2m$.

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The Dynkin diagram is \tilde{E}_8 :

6. A = (1, 1, 1, 1). Here r = 1, s = 4 and $n_1 = n_2 = n_3 = n_4 = m$, $m_1 = 2m$. The corresponding Dynkin diagram is \tilde{D}_4 .

7. $a_{1,1} = a_{1,2} = a_{1,3} = 1$, $a_{2,3} = a_{2,4} = a_{3,4} = a_{3,5} = \cdots = a_{k-2,k-1} = a_{k-2,k} = 1$, $a_{k-1,k} = a_{k-1,k+1} = a_{k-1,k+2} = 1$, and $a_{i,j} = 0$ for other (i, j).

Here we have r = k-1, s = k+2 and $n_1 = n_2 = n_{k+1} = n_{k+2} = m$, $n_3 = \cdots = n_k = 2m$, $m_1 = \cdots = m_l = 2m$. The corresponding Dynkin diagram is \tilde{D}_{2k} , where $k \ge 3$.



8.
$$a_{1,1} = a_{1,2} = a_{1,3} = 1$$
, $a_{2,3} = a_{2,4} = a_{3,4} = a_{3,5} = \cdots = a_{k-2,k-1} = a_{k-2,k} = 1$, $a_{k-1,k} = a_{k,k} = 1$, and $a_{i,j} = 0$ for other (i, j) .

Here we have $r = s = k \ge 3$, $n_1 = n_2 = m$, $n_3 = \cdots = n_k = 2m$, $m_1 = \cdots = m_{k-2} = 2m$, $m_{k-1} = m_k = m$. The corresponding Dynkin diagram is \tilde{D}_{2k-1} :



Note that if k = 3, then $a_{1,1} = a_{1,2} = a_{1,3} = 1$, $a_{2,3} = a_{3,3} = 1$.

Resume: Suppose \mathcal{L} is an indecomposable \mathcal{M} -structure with semi-simple algebras

 $\mathcal{A} = \bigoplus_{1 \leq i \leq r} End(V_i)$, $\mathcal{B} = \bigoplus_{1 \leq j \leq s} End(W_j)$; then there exists an affine Dynkin diagram of the type A, D, or E such that:

1. There is a one-to-one correspondence between the set of vertices and the set of vector spaces

 $\{V_1, ..., V_r, W_1, ..., W_s\}$

. **2.** For any i, j the spaces V_i , V_j are not connected by edges as well as W_i , W_j .

3. The vector

 $(\dim V_1, ..., \dim V_r, \dim W_1, ..., \dim W_s)$

is equal to mJ, where J is the minimal imaginary positive root of the Dynkin diagram.

Remark. It can be proved that for indecomposable \mathcal{M} -structures m = 1.

To describe the \mathcal{M} -structure it remains to construct an embedding $\mathcal{A} \to \mathcal{L}$, $\mathcal{B} \to \mathcal{L}$ and a scalar product (\cdot, \cdot) on the space \mathcal{L} .

If we fix an element $1 \in \mathcal{L}$, then we can define the embedding $\mathcal{A} \to \mathcal{L}$, $\mathcal{B} \to \mathcal{L}$ by the formula $a \to 1a$, $b \to b1$ for $a \in \mathcal{A}$, $b \in \mathcal{B}$. We may assume that 1 is a generic element of \mathcal{L} .

Thus to study \mathcal{M} -structures corresponding to a Dynkin diagram, one should take a generic element in $\mathcal{L} = \bigoplus_{1 \leq i \leq r, 1 \leq j \leq s} (V_i^{\star} \otimes W_j)^{a_{i,j}}$, find its simplest canonical form by choosing bases in the vector spaces

$$V_1, ..., V_r, W_1, ..., W_s,$$

calculate the embedding $\mathcal{A} \to \mathcal{L}$, $\mathcal{B} \to \mathcal{L}$ and the scalar product (\cdot, \cdot) on the space \mathcal{L} .

The classification of generic elements $1 \in \mathcal{L}$ up to choice of bases in the vector spaces $V_1, ..., V_r, W_1, ..., W_s$ is equivalent to classification of irreducible representations of the quivers corresponding to our affine Dynkin diagrams and we can apply known results about these representations.

Semi-simple case.

Consider an associative algebra

$$M = \oplus_{1 \le \alpha \le m} M_{\alpha},$$

where M_{α} is isomorphic to $Mat_{n_{\alpha}}$. We are going to study associative products in this algebra compatible with the initial one.

The main algebraic object is the weak \mathcal{PM} -structure (of size m) on \mathcal{L} . All properties coincide with the properties of weak \mathcal{M} -structure except for

1. dim $\mathcal{A} \cap \mathcal{B} = \dim \mathcal{L}/(\mathcal{A} + \mathcal{B}) = m$. The intersection of \mathcal{A} and \mathcal{B} is a *m*-dimensional algebra isomorphic to $\mathbb{C} \oplus ... \oplus \mathbb{C}$.

The scalar product (\cdot, \cdot) defines a non- degenerate pairing between $\mathcal{A}/\mathcal{A} \cap \mathcal{B}$ and $\mathcal{B}/\mathcal{A} \cap \mathcal{B}$, so dim $\mathcal{A} = \dim \mathcal{B}$ and dim $\mathcal{L} = 2 \dim \mathcal{A}$.

Let $\{e_{\alpha}; 1 \leq \alpha \leq m\}$ be a basis of the space $\mathcal{A} \cap \mathcal{B}$ such that

$$e_{\alpha}e_{\beta} = \delta_{\alpha,\beta}e_{\alpha}.$$

Denote by $\mathcal{L}_{\alpha,\beta}$ the vector space consisting of elements $v_{\alpha,\beta} \in \mathcal{L}$ with the property

$$e_{\alpha}v_{\alpha,\beta} = v_{\alpha,\beta}e_{\beta} = v_{\alpha,\beta}.$$

Let $\mathcal{A}_{\alpha,\beta} = \mathcal{A} \cap \mathcal{L}_{\alpha,\beta}$ and $\mathcal{B}_{\alpha,\beta} = \mathcal{B} \cap \mathcal{L}_{\alpha,\beta}.$

We sum by repeated Latin indexes and do not sum by repeated Greek indexes

The following properties hold:

- $\mathcal{L} = \bigoplus_{1 \leq \alpha, \beta \leq m} \mathcal{L}_{\alpha, \beta}$, $\mathcal{A} = \bigoplus_{1 \leq \alpha, \beta \leq m} \mathcal{A}_{\alpha, \beta}$ and $\mathcal{B} = \bigoplus_{1 \leq \alpha, \beta \leq m} \mathcal{B}_{\alpha, \beta}$.
- dim $\mathcal{A}_{\alpha,\beta} \cap \mathcal{B}_{\alpha,\beta} = \dim \mathcal{L}/(\mathcal{A}_{\alpha,\beta} + \mathcal{B}_{\alpha,\beta}) = \delta_{\alpha,\beta}$. The intersection of $\mathcal{A}_{\alpha,\alpha}$ and $\mathcal{B}_{\alpha,\alpha}$ is an one-dimensional space spanned by e_{α} .
- $\mathcal{B}_{\alpha,\beta}\mathcal{L}_{\beta',\gamma} = 0$ for $\beta \neq \beta'$ and $\mathcal{B}_{\alpha,\beta}\mathcal{L}_{\beta,\gamma} \subset \mathcal{L}_{\alpha,\gamma}$. Similarly $\mathcal{L}_{\alpha,\beta}\mathcal{A}_{\beta',\gamma} = 0$ for $\beta \neq \beta'$ and $\mathcal{L}_{\alpha,\beta}\mathcal{A}_{\beta,\gamma} \subset \mathcal{L}_{\alpha,\gamma}$. In particular, $\mathcal{A}_{\alpha,\beta}\mathcal{A}_{\beta',\gamma} = \mathcal{B}_{\alpha,\beta}\mathcal{B}_{\beta',\gamma} = 0$ for $\beta \neq \beta'$ and $\mathcal{A}_{\alpha,\beta}\mathcal{A}_{\beta,\gamma} \subset \mathcal{A}_{\alpha,\gamma}$, $\mathcal{B}_{\alpha,\beta}\mathcal{B}_{\beta,\gamma} \subset \mathcal{B}_{\alpha,\gamma}$.

• $\mathcal{L}_{\alpha,\beta} \perp \mathcal{L}_{\beta',\alpha'}$ if $\alpha \neq \alpha'$ or $\beta \neq \beta'$.

It follows from these properties that (\cdot, \cdot) gives nondegenerate pairing between $\mathcal{A}_{\alpha,\beta}$ and $\mathcal{B}_{\beta,\alpha}$ for $\alpha \neq \beta$ and between $\mathcal{A}_{\alpha,\alpha}/\mathbb{C}e_{\alpha}$ and $\mathcal{B}_{\alpha,\alpha}/\mathbb{C}e_{\alpha}$. Therefore dim $\mathcal{A}_{\alpha,\beta} =$ dim $\mathcal{B}_{\beta,\alpha}$.

Definition. By weak \mathcal{PM} -algebra associated with a weak \mathcal{PM} -structure \mathcal{L} we mean an associative algebra $U(\mathcal{L})$ possessing the following properties:

1. $\mathcal{L} \subset U(\mathcal{L})$ and the actions $\mathcal{B} \times \mathcal{L} \to \mathcal{L}$, $\mathcal{L} \times \mathcal{A} \to \mathcal{L}$ are the restrictions of the product in $U(\mathcal{L})$.

2. For any algebra X with the property **1** there exists a unique homomorphism of algebras $X \to U(\mathcal{L})$ identical on \mathcal{L} .

It is clear that $U(\mathcal{L}) = \bigoplus_{1 \leq \alpha, \beta \leq m} U(\mathcal{L})_{\alpha,\beta}$, where $U(\mathcal{L})_{\alpha,\beta} = \{v \in U(\mathcal{L}); e_{\alpha}v = ve_{\beta} = v\}$. We have $U(\mathcal{L})_{\alpha,\beta}U(\mathcal{L})_{\beta',\gamma} = 0$ for $\beta \neq \beta'$ and $U(\mathcal{L})_{\alpha,\beta}U(\mathcal{L})_{\beta,\gamma} \subset U(\mathcal{L})_{\alpha,\gamma}$.

Let $\{e_{\alpha}, A_{i,\alpha,\alpha}; 1 \leq i \leq p_{\alpha,\alpha}\}$ be a basis of $\mathcal{A}_{\alpha,\alpha}$ and $\{e_{\alpha}, B^{i}_{\alpha,\alpha}; 1 \leq i \leq p_{\alpha,\alpha}\}$ be the dual basis of $\mathcal{B}_{\alpha,\alpha}$. Let $\{A_{i,\alpha,\beta}; 1 \leq i \leq p_{\beta,\alpha}\}$ be a basis of $A_{\alpha,\beta}$ for $\alpha \neq \beta$ and $\{B^{i}_{\beta,\alpha}; 1 \leq i \leq p_{\beta,\alpha}\}$ be the dual basis of $B_{\beta,\alpha}$.

Take $C_{\alpha} \in \mathcal{L}_{\alpha,\alpha}$ that does not belong to the sum of $\mathcal{A}_{\alpha,\alpha}$ and $\mathcal{B}_{\alpha,\alpha}$. Without loss of generality we can assume that $(e_{\alpha}, C_{\alpha}) = 1$, $(C_{\alpha}, C_{\alpha}) = (C_{\alpha}, A_{i,\alpha,\alpha}) = (C_{\alpha}, B^{j}_{\alpha,\alpha}) = 0$. Such element C_{α} is uniquely determined by choosing of basis in $\mathcal{A}_{\alpha,\alpha}$.

Let us define the element $K_{\alpha} \in U(\mathcal{L})$ by the formula $K_{\alpha} = C_{\alpha} + \sum_{1 \leq \nu \leq m} A_{i,\alpha,\nu} B^{i}_{\nu,\alpha}$.

Definition. A weak \mathcal{PM} -structure \mathcal{L} is called \mathcal{PM} structure if $K = \sum_{1 \le \alpha \le m} K_{\alpha} \in U(\mathcal{L})$ is a central element of the algebra $U(\mathcal{L})$. **Theorem 3.4.** A basis of $U(\mathcal{L})_{\alpha,\beta}$ for $\alpha \neq \beta$ consists of the elements

$$\{A_{i,\alpha,\beta}K_{\beta}^{s}, B_{\alpha,\beta}^{j}K_{\beta}^{s}, A_{i_{1},\alpha,\nu}B_{\nu,\beta}^{j_{1}}K_{\beta}^{s}\},$$
where $1 \leq i \leq p_{\beta,\alpha}, 1 \leq j \leq p_{\alpha,\beta}, 1 \leq \alpha, \beta, \nu \leq m, 1 \leq i_{1} \leq p_{\nu,\alpha}, 1 \leq j_{1} \leq p_{\nu,\beta}, s = 0, 1, 2, \dots$

A basis of $U(\mathcal{L})_{\alpha,\alpha}$ consists of the elements

 $\{e_{\alpha}, A_{i,\alpha,\alpha}K_{\alpha}^{s}, B_{\alpha,\alpha}^{j}K_{\alpha}^{s}, A_{i_{1},\alpha,\nu}B_{\nu,\alpha}^{j_{1}}K_{\alpha}^{s}\},$ where $1 \leq i, j \leq p_{\alpha,\alpha}, 1 \leq \nu \leq m, 1 \leq i_{1}, j_{1} \leq p_{\nu,\alpha}, s = 0, 1, 2, \dots$ **Theorem.** Let \mathcal{L} be a \mathcal{PM} -structure. Then for any representation of $U(\mathcal{L})$ given by

$$x_{\alpha} \circ y_{\beta} = a_{i,\beta,\alpha} x_{\alpha} b^{i}_{\alpha,\beta} y_{\beta} + x_{\alpha} a_{i,\alpha,\beta} y_{\beta} b^{i}_{\beta,\alpha}, \quad \alpha \neq \beta,$$

$$x_{\alpha} \circ y_{\alpha} = a_{i,\alpha,\alpha} x_{\alpha} b^{i}_{\alpha,\alpha} y_{\alpha} + x_{\alpha} a_{i,\alpha,\alpha} y_{\alpha} b^{i}_{\alpha,\alpha} -$$

$$a_{i,\alpha,\alpha}x_{\alpha}y_{\alpha}b_{\alpha,\alpha}^{i} + x_{\alpha}c_{\alpha}y_{\alpha}.$$

defines a product on $M = \bigoplus_{1 \le \alpha \le m} Mat_{n_{\alpha}}$ compatible with the usual one.

Lemma 2. If (5) is associative, then for some tensor t_j^i

$$\phi_{j,k}^{s}\psi_{s}^{l,i} = \phi_{s,k}^{l}\psi_{j}^{s,i} + \phi_{j,s}^{i}\psi_{k}^{l,s} + \delta_{k}^{l}t_{j}^{i} - \delta_{j}^{i}t_{k}^{l} - \delta_{j}^{i}\phi_{s,r}^{l}\psi_{k}^{r,s},$$
 and

$$b^{i}a_{j} = \psi_{j}^{k,i} a_{k} + \phi_{j,k}^{i} b^{k} + t_{j}^{i} + \delta_{j}^{i} c.$$

Lemma 3. If (5) is associative, then

$$b^{i} c = \lambda^{k,i} a_{k} - t_{k}^{i} b^{k} - \phi_{k,l}^{i} \psi_{s}^{l,k} b^{s} - \phi_{k,l}^{i} \lambda^{l,k},$$

$$c a_j = \mu_{j,k} b^k - t_j^k a_k - \phi_{k,l}^s \psi_j^{l,k} a_s - \mu_{k,l} \psi_j^{l,k}$$

and moreover

$$\phi_{j,k}^{s}t_{s}^{i} = \psi_{j}^{s,i}\mu_{s,k} + \phi_{j,s}^{i}t_{k}^{s} - \delta_{j}^{i}\psi_{k}^{s,r}\mu_{r,s},$$
$$\psi_{s}^{k,i}t_{j}^{s} = \phi_{j,s}^{i}\lambda^{k,s} + \psi_{j}^{s,i}t_{s}^{k} - \delta_{j}^{i}\phi_{s,r}^{k}\lambda^{r,s}.$$

Proposition 1. The algebra $U(\mathcal{L})$ is defined by the following relations

$$\begin{split} A_i A_j &= \phi_{i,j}^k A_k + \mu_{i,j}, \qquad B^i B^j = \psi_k^{i,j} B^k + \lambda^{i,j} \\ B^i A_j &= \psi_j^{k,i} A_k + \phi_{j,k}^i B^k + t_j^i + \delta_j^i C, \\ B^i C &= \lambda^{k,i} A_k + u_k^i B^k + p^i, \quad C A_j = \mu_{j,k} B^k + u_j^k A_k + q_i \\ \text{for certain tensors } \phi_{i,j}^k, \psi_k^{i,j}, \mu_{i,j}, \lambda^{i,j}, u_k^i, p^i, q_i. \end{split}$$

Lemma. \mathcal{L} as \mathcal{A} -module is isomorphic to $\bigoplus_{1 \le i \le r} (V_i^{\star})^{2m_i}$.

Definition. The $r \times s$ matrix $A = (a_{i,j})$ is called matrix of multiplicities of a weak \mathcal{M} -structure.

Definition. The matrix A is called decomposable if there exist partitions $\{1, ..., r\} = I \sqcup I'$ and $\{1, ..., s\} = J \sqcup J'$ such that $a_{i,j} = 0$ for $(i, j) \in I \times J' \sqcup I' \times J$.

Lemma. If A is decomposable, then the corresponding \mathcal{M} -structure is decomposable.